

The logistic Keller–Segel equation and its particle approximation

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Keller–Segel model: Motivation

- ▶ System of PDE describing the collective motion of cells (typically bacteria), attracted by a chemical substance (chemoattractant) that they themselves emit.
- ▶ The parabolic-elliptic form is

$$\begin{cases} \partial_t \rho = \Delta \rho - \chi \nabla \cdot (\rho \nabla c), & t > 0, x \in \mathbb{R}^d \\ -\Delta c = \rho, & t > 0, x \in \mathbb{R}^d, \end{cases}$$

where $\chi > 0$ is the intensity of chemoattractant.

- ▶ Solving Poisson's equation for c gives $\nabla c = K * \rho$, where $K(x) \propto -\frac{x}{|x|^d}$, and so

$$\partial_t \rho = \Delta \rho - \chi \nabla \cdot (\rho K * \rho).$$

- ▶ Competition between diffusion and attraction.
- ▶ For large χ solutions blow up in finite time (mass concentration), e.g. $\chi > \frac{8\pi}{m}$ in $d = 2$ [Per04].

Keller–Segel with logistic source

- ▶ There are several mechanisms to suppress or delay blow up, one of which is a logistic source term modelling births and deaths.
- ▶ Introduced by Tello and Winkler [TW07] on bounded domains and recently extended to $\mathbb{R}^d$ by Cavallazzi, Richard and Tomašević [CRT25]:

$$\partial_t \rho = \Delta \rho - \chi \nabla \cdot (\rho K * \rho) + \nu \rho - \mu \rho^2,$$

where $\nu \geq 0$ is the birth rate and $\mu \rho \geq 0$ is the death rate.

[CRT25, Thm. 1] (Well-posedness)

For $d \geq 2$, for any $\nu \geq 0$, if

$$\mu > \frac{d-2}{d} \chi,$$

there exist a unique global mild solution.

Quantitative approximation by particles I

- ▶ A central question is whether the logistic Keller–Segel equation can be derived from a system of interacting particles.
- ▶ We consider a logistic particle system starting with N particles using a Poisson process:
 - ▶ particles branch at rate ν ;
 - ▶ particles die at a rate proportional to the local density.
- ▶ Between birth and death events, particles evolve according to

$$dX_t^k = \chi F_A \left(\frac{1}{N} \sum_{j \in I_t} K * \theta^N(X_t^k - X_t^j) \right) dt + \sqrt{2} dB_t^k,$$

where $I_t =$ particles alive at time t .

- ▶ The interaction is regularised by a cutoff F_A and mollifier θ^N , Oelschläger [Oel85]. Used for technical reasons.

Quantitative approximation by particles II

- ▶ Let $\mu_t^N := \frac{1}{N} \sum_{k \in I_t} \delta_{X_t^k}$ denote the empirical measure and $\rho_t^N := \theta^N * \mu_t^N$ its mollification.

[CRT25, Thm. 2] (Mean field limit)

Whenever $\mu > \frac{d-2}{d} \chi$, for any $T, \epsilon > 0$ and initial number of particles N , there exists an explicit rate $\Gamma > 0$ such that

$$\sup_{t \in [0, T]} \|\rho_t^N - \rho_t\| \lesssim \|\rho_0^N - \rho_0\| + N^{-\Gamma + \epsilon}.$$

- ▶ “Large population” result.
- ▶ The same convergence rate holds for the empirical measure itself (in Kantorovich–Rubinstein distance).

Current work: Removing moderate interactions

- ▶ Remark: moderate interacting particles were used because the well-posedness of the singular system particle is not known.
- ▶ In dimension $d = 2$, the classical supercritical Keller–Segel interacting particle system $\chi > \frac{8\pi}{m}$ around the time of blow up was analysed by Fournier and Tardy [FT25]

$$dX_t^k = \frac{\chi}{N} \sum_{j \neq k} \frac{X_t^j - X_t^k}{\|X_t^j - X_t^k\|^2} dt + \sqrt{2} dB_t^k,$$

- ▶ They show that at the time of explosion, there is a critical proportion $k_0 = \lceil N \frac{8\pi}{m\chi} \rceil$ of particles which collide together.
- ▶ Our aim is to analyse the logistic particle system, and show that the killing mechanism prevents there from being a critical proportion of particles close together.

Current results: The particle system I

- ▶ A difference in the microscopic and macroscopic system.
- ▶ No blow up of the PDE, but it is possible with low probability that the killing mechanism does not kill aggregating particles.

Theorem (No critical collision with high prob.)

Let $d = 2$, fix terminal time $T > 0$ and arbitrary killing rate $\mu > 0$.

For any $\epsilon > 0$, there exists an initial number of particles N sufficiently large such that

$\mathbb{P}(\text{particle system has a critical collision before time } T) < \epsilon.$

Current results: proof of no critical collisions

Proof idea

- ▶ Introduce a functional $F^\gamma(\mu_t^N)$ measuring the proportion of particles that are γ distance apart.
- ▶ Itô formula on the particles tells us that if proportion exceeds level $k_0/2$, the killing mechanism dominates and generates a strong negative drift:

$$F_t^\gamma \lesssim \frac{k_0}{2} + \frac{1}{N}M_t - Nc_{\nu,\mu},$$

with $c_{\nu,\mu} > 0$ due to the killing mechanism.

- ▶ The negative drift overwhelms the martingale fluctuations as $N \rightarrow \infty$, preventing F^γ from reaching the critical threshold k_0 with high probability.

Current results II: Towards a complete mean field limit

Theorem (Mean field limit, $d = 2$)

- ▶ The empirical measures $((\mu_t^N)_{t \in [0, T]})_N$ are tight.
- ▶ Every sub-sequential limit is a “very weak” solution of the logistic Keller–Segel equation, in the sense of Tardy [Tar24].

Next steps

- ▶ Establish uniqueness of the limiting law.
 - ▶ Show that the limiting measure admits a density.
 - ▶ Upgrade the limit to a mild solution of the PDE.
 - ▶ Invoke uniqueness of mild solutions [CRT25].
- ▶ Upgrade subsequential convergence to convergence of the full sequence.
- ▶ Establish stronger modes of convergence (e.g. pathwise/trajectorial convergence).

Thank you for your attention, Obrigado :)

Any questions?

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