

# Strong periodic solutions for a class of abstract evolution equations

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Abstract. We consider a class of abstract evolution equations in a Hilbert space for which we prove existence and uniqueness of strong time periodic solutions. The result covers many models in hydrodynamics.

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# 1. Introduction.

We prove existence and uniqueness of strong periodic solutions for a class of abstract evolution equations of the form

$$(1.1) \quad \frac{d}{dt}U + \mathcal{L}U + B(U, U) + RU = F$$

in a Hilbert space  $\mathcal{H}$ . As an application we consider the system of equations of magneto-micropolar fluids. Our result covers several other models in hydrodynamics, including, e.g., Navier-Stokes equations with nonhomogenous boundary conditions, magnetohydrodynamics equations, and micropolar fluid equations.

Our assumptions are as follows.

$\mathcal{L}$  is a self-adjoint, strictly positive operator in  $\mathcal{H}$  with domain  $\mathcal{D}(\mathcal{L})$  and compact inverse.

$B$  is a bilinear form in  $\mathcal{H} \times \mathcal{H}$  such that

$$(1.2) \quad |\mathcal{B}(U, V, W)| \leq c_1 |\mathcal{L}^\theta U| \cdot |\mathcal{L}^\rho U| \cdot |W|$$

for  $U \in D(\mathcal{L}^\theta)$ ,  $V \in D(\mathcal{L}^\rho)$ ,  $W \in \mathcal{H}$ , where  $\mathcal{B}(U, V, W) = (B(U, V), W)$ ,  $(\cdot, \cdot)$  and  $|\cdot|$  denote scalar product and norm in  $\mathcal{H}$ , respectively,

$$(1.3) \quad \theta + \rho \geq \frac{5}{4}, \quad \rho > \frac{1}{2}, \quad \theta > 0,$$

and

$$(1.4) \quad (B(U, V), V) = 0, \quad U, V \in D(\mathcal{L}^{1/2})$$

$R$  is a linear operator in  $\mathcal{H}$  such that

$$(1.5) \quad |\mathcal{R}(U, V)| \leq c_2 |\mathcal{L}^{1/2}U| \cdot |V|$$

where  $\mathcal{R}(U, V) = (RU, V)$ , and

$$(1.6) \quad \mathcal{A}(U, U) + \lambda \mathcal{R}(U, U) \geq k_1 |\mathcal{L}^{1/2}U|^2, \quad \lambda \in [0, 1]$$

where  $\mathcal{A}(U, V) = (\mathcal{L}^{1/2}U, \mathcal{L}^{1/2}V)$ ,  $U, V \in D(\mathcal{L}^{1/2})$ ,  $k_1 > 0$ . We prove the following

**Theorem 1.1** *Let the assumptions (1.2)-(1.6) hold, and let  $F \in C^1(\tau, \mathcal{H})$  for some  $\tau > 0$ . Then there exists a constant  $M$  such that if*

$$(1.7) \quad \sup_t |F(t)| \leq M$$

then (1.1) has a unique strong  $\tau$ -periodic solution

$$(1.8) \quad U \in H^2(\tau; \mathcal{H}) \cap H^1(\tau; D(\mathcal{L})) \cap L^\infty(\tau; D(\mathcal{L})) \cap W^{1,\infty}(\tau; D(\mathcal{L}^{1/2})) \quad \blacksquare$$

By  $C^k(\tau, \mathcal{H})$ ,  $k = 0, 1, 2, \dots$ , we denote the Banach space of  $\mathcal{H}$ -valued  $\tau$ -periodic functions on  $R^1$  with continuous derivatives up to order  $k$ , with the usual norm

$$\|f\|_{C^k(\tau, \mathcal{H})} = \sup \left\{ \sum_{i=0}^k |D_t^i f(t)| : 0 \leq t \leq \tau \right\}$$

For a Hilbert space  $X$ , by  $H^k(\tau, X)$ ,  $k = 0, 1, 2, \dots$ , we denote the Hilbert space of  $X$ -valued functions on  $R^1$  that belong, together with their derivatives up to order  $k$ , to the space of  $\tau$ -periodic measurable functions  $L^2(\tau; X)$  with the norm  $\|f\|_{L^2(\tau; X)} = (\int_0^\tau \|f(t)\|_X^2 dt)^{1/2}$ . Similarly, we define Banach spaces  $L^\infty(\tau; X)$  and  $W^{1,\infty}(\tau; X)$  with norms

$$\|f\|_{L^\infty(\tau; X)} = \sup \text{ess} \{ \|f(t)\|_X : 0 \leq t \leq \tau \}$$

and

$$\|f\|_{W^{1,\infty}(\tau; X)} = (\|f\|_{L^\infty(\tau; X)}^2 + \|D_t f\|_{L^\infty(\tau; X)}^2)^{1/2}$$

By  $C$  we denote various numeric constants.

The proof of theorem 1.1 is based on a set of estimates applied to the family  $U_n \in C^1(\tau; \mathcal{H}_n)$ ,  $n = 1, 2, 3, \dots$  of approximate solutions, where  $\mathcal{H}_n = \text{lin}\{\omega_1, \omega_2, \dots, \omega_n\}$  and  $(\omega_i)$  is a complete orthonormal system in  $\mathcal{H}$  consisting of eigenfunctions of the operator  $\mathcal{L}$ , cf. [5].

In section 2 we prove the existence of solutions  $U_n$ . In section 3 we obtain a set of estimates of  $U_n$  that allows us to pass to the limit with  $n$ . In section 4 we prove the existence of a solution of equation (1.1) satisfying (1.8) by letting  $n$  to infinity as well as its uniqueness.

Section 5 presents an application of theorem 1.1 to the system of equations of magneto-micropolar fluids.

## 2. Approximate solutions and first order estimates.

Let

$$(2.1) \quad (U_t, \omega_i) + \mathcal{A}(U, \omega_i) + \mathcal{B}(U, U, \omega_i) + \mathcal{R}(U, \omega_i) = (F, \omega_i), \quad i = 1, 2, \dots, n$$

where  $U(t) = U_n(t) = \sum_{i=1}^n c_{in}(t)\omega_i$ ,  $U(t) = U(t + \tau)$ ,  $n = 1, 2, 3, \dots$  be the system of ODE defining approximate solutions of equation (1.1).

**Lemma 2.1** *For each  $n$  there exists a solution  $U \in C^1(\tau, \mathcal{H}_n)$  of the nonlinear problem (2.1).*

Proof. The linear problem

$$(2.2) \quad (U_t, \phi) + \mathcal{A}(U, \phi) = (F, \phi) - \mathcal{R}(V, \phi) - \mathcal{B}(V, V, \phi), \quad \phi \in \mathcal{H}_n,$$

$$U(t) = U(t + \tau)$$

has a unique solution (cf. [1], [2])  $U \in C^1(\tau, \mathcal{H}_n)$  for each  $V \in C^0(\tau, \mathcal{H}_n)$ . Consider the map  $\Phi : V \rightarrow U$  in the space  $C^0(\tau, \mathcal{H}_n)$ . We shall show that  $\Phi$  has a fixed point by using the Leray-Schauder theorem.

We prove that for every  $U$  and  $\lambda \in [0, 1]$  satisfying  $\lambda\Phi(U) = U$ ,

$$(2.3) \quad \sup_{0 \leq t \leq \tau} |U(t)| \leq CM \quad \text{for some } C > 0$$

where  $M$  is a constant in (1.7).

For  $\lambda = 0$ ,  $U = 0$ . Let  $\lambda > 0$  and assume  $\lambda\Phi(U) = U$ . Then, from (2.2)

$$(U_t, \phi) + \mathcal{A}(U, \phi) = \lambda(F, \phi) - \lambda\mathcal{R}(U, \phi) - \lambda\mathcal{B}(U, U, \phi)$$

With  $\phi = U$ , in view of (1.4), we obtain

$$\frac{1}{2} \frac{d}{dt} |U|^2 + \mathcal{A}(U, U) + \lambda\mathcal{R}(U, U) = \lambda(F, U)$$

Denoting by  $\mu$  the smallest eigenvalue of  $\mathcal{L}$  we have

$$(2.4) \quad |\mathcal{L}^\alpha U| \leq \mu^{\alpha-\beta} |\mathcal{L}^\beta U|, \quad 0 \leq \alpha \leq \beta,$$

whence  $|(F, U)| \leq C|F| |\mathcal{L}^{1/2} U|$ , and by (1.6) we obtain

$$(2.5) \quad \frac{d}{dt} |U|^2 + k_1 |\mathcal{L}^{1/2} U|^2 \leq CM^2$$

Integrating in  $t$  and using the periodicity of  $U$  we have

$$\int_0^\tau |\mathcal{L}^{1/2} U(t)|^2 dt \leq CM^2 \tau,$$

whence, by the mean value theorem for integrals and (2.4), there exists  $t^* \in [0, \tau]$  such that

$$(2.6) \quad |U(t^*)| \leq C|\mathcal{L}^{1/2}U(t^*)| \leq CM$$

Integrating again (2.5) from  $t^*$  to  $t + \tau$ ,  $t \in [0, \tau]$  we obtain (2.3). As the map  $\Phi$  is continuous and compact in  $C^0(\tau; \mathcal{H}_n)$  we conclude the existence of a fixed point  $U$  for  $\Phi$ . Observe that (2.3) holds for this  $U$ . ■

Now we shall prove a fundamental estimate.

**Lemma 2.2** *Let  $U = U_n$  be the solution from (2.1). Then there exists a constant  $C$  independent of  $n$  and such that*

$$(2.7) \quad \sup_{t \in \mathbb{R}^1} |\mathcal{L}^{1/2}U(t)| \leq CM^{1/2}$$

Proof. From (2.6) with  $M < 1$  we have

$$|\mathcal{L}^{1/2}U(t^*)| < CM^{1/2}$$

Let  $T^* = \sup\{T : |\mathcal{L}^{1/2}U(t)| \leq CM^{1/2} \text{ for } t \in [t^*, T]\}$ . We shall show that  $T^* = \infty$ , whence (2.7) by periodicity. Let us assume to the contrary that  $T^* < \infty$ . Then

$$(2.8) \quad |\mathcal{L}^{1/2}U(T^*)| = CM^{1/2}$$

Our aim is to derive, using (2.8), the differential inequality

$$(2.9) \quad \frac{d}{dt} |\mathcal{L}^{1/2}U(T^*)|^2 \leq 0$$

that shows the contradiction in view of the definition of  $T^*$ .

For all  $\phi \in \mathcal{H}_n$  we have

$$(2.10) \quad (U_t, \phi) + \mathcal{A}(U, \phi) + \mathcal{B}(U, U, \phi) + \mathcal{R}(U, \phi) = (F, \phi)$$

Setting  $\phi = \mathcal{L}U$  we obtain

$$\frac{1}{2} \frac{d}{dt} |\mathcal{L}^{1/2}U|^2 + |\mathcal{L}U|^2 \leq |F| |\mathcal{L}U| + |\mathcal{R}(U, \mathcal{L}U)| + |\mathcal{B}(U, U, \mathcal{L}U)|$$

where we have used  $(\mathcal{L}U, \phi) = \mathcal{A}(U, \phi)$ ,  $U \in D(\mathcal{L})$ ,  $\phi \in \mathcal{H}$ . Now by (1.2) and (1.5),

$$(2.11) \quad \frac{1}{2} \frac{d}{dt} |\mathcal{L}^{1/2}U|^2 + |\mathcal{L}U|^2 \leq |F| |\mathcal{L}U| + c_2 |\mathcal{L}^{1/2}U| \cdot |\mathcal{L}U| + c_1 |\mathcal{L}^{1/4}U| \cdot |\mathcal{L}U|^2$$

Let us consider this inequality at point  $t = T^*$ . Then, by (2.3), (2.8), and the interpolation inequality

$$(2.12) \quad |\mathcal{L}^\sigma U| \leq c_2 |\mathcal{L}^\alpha U|^\gamma \cdot |\mathcal{L}^\beta U|^{1-\gamma}, \quad U \in D(\mathcal{L}^\beta)$$

where  $\sigma = \alpha\lambda + \beta(1 - \lambda)$ ,  $0 \leq \alpha \leq \sigma \leq \beta$ ,  $\lambda \geq 0$ , we have

$$(2.13) \quad |\mathcal{L}^{1/4} U| \cdot |\mathcal{L}U|^2 \leq C|U|^{1/2} \cdot |\mathcal{L}^{1/2} U|^{1/2} \cdot |\mathcal{L}U|^2 \leq CM^{1/2} M^{1/4} |\mathcal{L}U|^2$$

$$(2.14) \quad \begin{aligned} |\mathcal{L}^{1/2} U| \cdot |\mathcal{L}U| &\leq \varepsilon |\mathcal{L}U|^2 + c_\varepsilon |\mathcal{L}^{1/2} U|^2 \\ &\leq \varepsilon |\mathcal{L}U|^2 + c_\varepsilon |U| \cdot |\mathcal{L}U| \\ &\leq \varepsilon |\mathcal{L}U|^2 + c_\varepsilon M \cdot |\mathcal{L}U| \\ &\leq \varepsilon |\mathcal{L}U|^2 + c_\varepsilon M^{1/2} \cdot |\mathcal{L}^{1/2} U| \cdot |\mathcal{L}U| \\ &\leq \varepsilon |\mathcal{L}U|^2 + c_\varepsilon M^{1/2} |\mathcal{L}U|^2 \end{aligned}$$

and, by (1.7),

$$(2.15) \quad |F||U| \leq CM|\mathcal{L}U| = CM^{1/2} |\mathcal{L}^{1/2} U| \cdot |\mathcal{L}U| \leq CM^{1/2} |\mathcal{L}U|^2$$

Thus, from (2.11), together with (2.13)-(2.15), we get

$$\frac{1}{2} \frac{d}{dt} |\mathcal{L}^{1/2} U|^2 + (1 - C(M)) |\mathcal{L}U|^2 \leq 0$$

where  $C(M) \searrow 0$  as  $M \searrow 0$ , whence (2.9), provided  $M$  is small enough. Q.E.D. ■

Now using (2.7) and (2.11) we have, for small  $M$ ,

$$\frac{d}{dt} |\mathcal{L}^{1/2} U|^2 + |\mathcal{L}U|^2 \leq |F|^2 + CM$$

In view of the periodicity of  $U$ , by integration we obtain

$$(2.16) \quad \int_0^\tau |\mathcal{L}U(t)|^2 dt \leq C(M, M_0), \quad M_0 \equiv \left( \int_0^\tau |F(t)|^2 dt \right)^{1/2}$$

**Lemma 2.3** *Let  $U$  be the solution from (2.1). Then*

$$(2.17) \quad \sup_{t \in \mathbb{R}^1} |U_t(t)|^2 \leq C(M, M_0, M_1), \quad C(M, M_0, M_1) \text{ independent of } n,$$

where  $M_0$  is as in (2.16) and  $M_1 = \left( \int_0^\tau |F_t(t)|^2 dt \right)^{1/2}$ .

Proof. Set  $\phi = U_t$  in (2.10). Then we have, by (1.5),

$$(2.18) \quad |U_t|^2 + \frac{1}{2} \frac{d}{dt} |\mathcal{L}^{1/2} U|^2 \leq |F| |U_t| + c_2 |\mathcal{L}^{1/2} U| \cdot |U_t| + |\mathcal{B}(U, U, U_t)|$$

In view of (1.2) and (2.7),

$$(2.19) \quad |\mathcal{B}(U, U, U_t)| \leq c_1 |\mathcal{L}^{1/4} U| \cdot |\mathcal{L} U| \cdot |U_t| \\ \leq C M^{1/2} |\mathcal{L} U|^2 + C M^{1/2} |U_t|^2$$

and from (2.18), (2.19), for small M,

$$|U_t|^2 + \frac{d}{dt} |\mathcal{L}^{1/2} U|^2 \leq |F|^2 + C |\mathcal{L}^{1/2} U|^2 + C M^{1/2} |\mathcal{L} U|^2$$

After integration in  $t$  we obtain, in view of (2.16), (2.7),

$$(2.20) \quad \int_0^\tau |U_t(t)|^2 dt \leq C(M, M_0).$$

Now, let us differentiate (2.10) with respect to  $t$  and set  $\phi = U_t$ . Then we get, by (1.4),

$$(U_{tt}, U_t) + \mathcal{A}(U_t, U_t) + \mathcal{B}(U_t, U, U_t) + \mathcal{R}(U_t, U_t) = (F_t, U_t)$$

whence, from (1.6),

$$\frac{d}{dt} |U_t|^2 + 2k_1 |\mathcal{L}^{1/2} U_t|^2 \leq C |F_t|^2 + \frac{k_1}{2} |\mathcal{L}^{1/2} U_t|^2 + 2|\mathcal{B}(U_t, U, U_t)|$$

By (1.2),

$$2|\mathcal{B}(U_t, U, U_t)| \leq 2c_1 |\mathcal{L}^{1/2} U_t| \cdot |\mathcal{L} U| \cdot |U_t| \\ \leq \frac{k_1}{2} |\mathcal{L}^{1/2} U_t|^2 + C |\mathcal{L} U|^2 \cdot |U_t|^2$$

From the two last inequalities we conclude

$$(2.21) \quad \frac{d}{dt} |U_t|^2 + k_1 |\mathcal{L}^{1/2} U_t|^2 \leq C |F_t|^2 + C |\mathcal{L} U|^2 \cdot |U_t|^2$$

and, in particular,

$$(2.22) \quad \frac{d}{dt} |U_t|^2 \leq C |F_t|^2 + C |\mathcal{L} U|^2 \cdot |U_t|^2$$

We shall use now the uniform Gronwall lemma [8]. We have

$$\int_t^{t+\tau} |U_t(s)|^2 ds = \int_0^\tau |U_t(s)|^2 ds \leq C(M, M_0)$$

by (2.20), and by (2.16)

$$\int_t^{t+\tau} |\mathcal{L} U(s)|^2 ds = \int_0^\tau |\mathcal{L} U(s)|^2 ds \leq C(M, M_0)$$

By our assumptions,

$$\int_t^{t+\tau} |F_t(s)|^2 ds = \int_0^\tau |F_t(s)|^2 ds \leq M_1^2$$

From the uniform Gronwall lemma applied to inequality (2.22) we obtain

$$|U_t(t + \tau)|^2 \leq \left\{ \frac{C(M, M_0)}{\tau} + M_1^2 \right\} \exp C(M, M_0) \quad \text{for all } t \geq 0$$

Since  $U$  is  $\tau$ -periodic, we obtain (2.17). Q.E.D. ■

### 3. Higher order estimates.

**Lemma 3.1** *Let  $U$  be the approximate solution from (2.1). Then*

$$(3.1) \quad \sup_{t \in \mathbb{R}^1} |\mathcal{L}U(t)| \leq C(M, M_0, M_1)$$

and

$$(3.2) \quad \sup_{t \in \mathbb{R}^1} |\mathcal{L}^{1/2}U_t(t)| \leq C(M, M_0, M_1)$$

Proof. Set  $\phi = \mathcal{L}U$  in (2.10), use (1.2) and (1.5), and then (2.17), (2.7), and (1.7) to get

$$\begin{aligned} |\mathcal{L}U|^2 &\leq |U_t| \cdot |\mathcal{L}U| + c_1 |\mathcal{L}^{1/2}U| \cdot |\mathcal{L}U| \cdot |\mathcal{L}U| + c_2 |\mathcal{L}^{1/2}U| \cdot |\mathcal{L}U| + |F| \cdot |\mathcal{L}U| \\ &\leq C(M, M_0, M_1) |\mathcal{L}U| + CM^{1/2} |\mathcal{L}U|^2 \end{aligned}$$

For  $M$  such that  $1 - CM^{1/2} > 0$  we obtain (3.1).

In order to obtain the second estimate of the lemma we differentiate identity (2.10) with respect to  $t$  and then set  $\phi = \mathcal{L}U_t$ . Using (1.2) and (1.5) we get

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} |\mathcal{L}^{1/2}U_t|^2 + |\mathcal{L}U_t|^2 &\leq c_1 |\mathcal{L}^{1/2}U_t| \cdot |\mathcal{L}^{3/4}U| \cdot |\mathcal{L}U_t| + c_1 |\mathcal{L}^{1/2}U| \cdot |\mathcal{L}^{3/4}U_t| \cdot |\mathcal{L}U_t| \\ &\quad + C |\mathcal{L}^{1/2}U_t|^2 + \frac{1}{4} |\mathcal{L}U_t|^2 + |F_t|^2 + \frac{1}{4} |\mathcal{L}U_t|^2 \end{aligned}$$

Using (2.4) and (3.1) we get

$$(3.3) \quad \frac{d}{dt} |\mathcal{L}^{1/2}U_t|^2 + |\mathcal{L}U_t|^2 \leq C(M, M_0, M_1) |\mathcal{L}^{1/2}U_t| \cdot |\mathcal{L}U_t|$$



$$+CM|\mathcal{L}U_t|^2 + C|\mathcal{L}^{1/2}U_t|^2 + |F_t|^2$$

From (2.21), (2.20),

$$(3.4) \quad \int_0^\tau |\mathcal{L}^{1/2}U_t|^2 dt \leq C(M, M_0, M_1)$$

and from (3.3), (3.4) and for  $M$  small enough,

$$\int_0^\tau |\mathcal{L}U_t|^2 dt \leq C(M, M_0, M_1)^{3/2} \cdot \left( \int_0^\tau |\mathcal{L}U_t|^2 dt \right)^{1/2} + C(M, M_0, M_1)$$

whence

$$(3.5) \quad \int_0^\tau |\mathcal{L}U_t(t)|^2 dt \leq C(M, M_0, M_1)$$

Thus, for some  $t^* \in [0, \tau]$ ,

$$|\mathcal{L}^{1/2}U_t(t^*)|^2 \leq C(M, M_0, M_1)$$

and, after integration of (3.3) in  $t$  from  $t^*$  to  $t + \tau$ ,  $t \in [0, \tau]$  we obtain (3.2). Q.E.D. ■

**Lemma 3.2** *Let  $U$  be the approximate solution from (2.1) Then*

$$(3.6) \quad \int_0^\tau |U_{tt}(t)|^2 dt \leq C(M, M_0, M_1)$$

Proof. Differentiate (2.10) with respect to  $t$  and set  $\phi = U_{tt}$ . We get

$$\begin{aligned} |U_{tt}|^2 &\leq |\mathcal{A}(U_t, U_{tt})| + |\mathcal{R}(U_t, U_{tt})| + |\mathcal{B}(U_t, U, U_{tt})| + |\mathcal{B}(U, U_t, U_{tt})| + |(F_t, U_{tt})| \\ &\leq |\mathcal{L}U_t| \cdot |U_{tt}| + C|\mathcal{L}^{1/2}U_t| \cdot |U_{tt}| + c_1|\mathcal{L}^{1/2}U_t| \cdot |\mathcal{L}U| \cdot |U_{tt}| + c_1|\mathcal{L}^{1/2}U| \cdot |\mathcal{L}U_t| \cdot |U_{tt}| + |F_t| \cdot |U_{tt}| \end{aligned}$$

in view of (1.2), (1.5), so that, by (2.4), (3.1), and (3.2),

$$|U_{tt}|^2 \leq C(M, M_0, M_1)|\mathcal{L}U_t|^2 + C|F_t|^2$$

Integration in  $t$  gives (3.6), by (3.5). Q.E.D. ■

## 4. Convergence and uniqueness.

From the estimates we have obtained it follows that  $(U_n)$  is a bounded sequence in

$$(4.1) \quad H^2(\tau; \mathcal{H}) \cap H^1(\tau; D(\mathcal{L})) \cap L^\infty(\tau; D(\mathcal{L})) \cap W^{1,\infty}(\tau; D(\mathcal{L}^{1/2}))$$

and that there exists a subsequence  $(U_\mu)$  and some  $U$  in the space (4.1) such that

$$U_\mu \rightarrow U \quad \text{weak star in } L^\infty(\tau; D(\mathcal{L}))$$

$$U_\mu \rightarrow U \quad \text{strongly in } L^\infty(\tau; D(\mathcal{L}^{1/2}))$$

$$D_t U_\mu \rightarrow D_t U \quad \text{weak star in } L^\infty(\tau; D(\mathcal{L}^{1/2}))$$

$$D_t U_\mu \rightarrow D_t U \quad \text{strongly in } L^\infty(\tau; \mathcal{H})$$

Set  $U = U_\mu$  in the identity (2.1) for  $\mu \geq n$ , and let  $\mu \rightarrow \infty$ . To obtain (1.1) observe that, by (1.2) and (2.12), for all  $\phi \in \mathcal{H}_n$  we have

$$|\mathcal{B}(U_\mu, U_\mu, \phi) - \mathcal{B}(U, U, \phi)| \leq |\mathcal{B}(U_\mu - U, U_\mu, \phi)| + |\mathcal{B}(U, U_\mu - U, \phi)|$$

$$\leq c_1 |\mathcal{L}^{1/2}(U_\mu - U)| \cdot |\mathcal{L}^{3/4} U_\mu| \cdot |\phi| + c_1 |\mathcal{L}^{1/2} U| \cdot |\mathcal{L}^{1/2}(U_\mu - U)|^{1/2} \cdot |\mathcal{L}(U_\mu - U)|^{1/2} |\phi| \rightarrow 0$$

uniformly in  $t$ , in view of (2.7), (3.1), and then

$$(U_t, \phi) + \mathcal{A}(U, \phi) + \mathcal{B}(U, U, \phi) + \mathcal{R}(U, \phi) = (F, \phi)$$

easily follows for the limit function  $U$  and all  $\phi \in \mathcal{H}$ .

To prove the uniqueness assume, to the contrary, that there are two different solutions  $U$  and  $V$ . Then  $W = U - V$  satisfies, for all  $\phi \in \mathcal{H}$ ,

$$(W_t, \phi) + \mathcal{A}(W, \phi) + \mathcal{B}(U, W, \phi) + \mathcal{B}(W, V, \phi) + \mathcal{R}(W, \phi) = 0$$

Set  $\phi = W$  and we get by (1.2), (2.12), (2.4) and (2.7), (3.1),

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} |W|^2 + k_1 |\mathcal{L}^{1/2} U|^2 &\leq |\mathcal{B}(W, V, W)| \leq c_1 |\mathcal{L}^{1/2} W| \cdot |\mathcal{L}^{3/4} V| \cdot |W| \leq \\ &\leq c_1 |\mathcal{L}^{1/2} W| \cdot |\mathcal{L}^{1/2} V|^{1/2} \cdot |\mathcal{L} V|^{1/2} \cdot |W| \leq M^{1/4} \cdot C(M, M_0, M_1) |\mathcal{L}^{1/2} W|^2 \end{aligned}$$

whence for small  $M$  we obtain, for some  $c > 0$ ,

$$(4.2) \quad \frac{d}{dt} |W|^2 + c |W|^2 \leq 0$$

which leads to the contradiction, as (4.2) yields  $W \equiv 0$  by Gronwall's lemma and periodicity of  $W$ . This proves the uniqueness. Q.E.D. ■

**Remark 4.1** Observe that in fact we have used (1.2)-(1.3) only for three pairs of parameters  $(\theta, \rho)$ , namely,  $(1/4, 1)$ ,  $(1/2, 1)$ , and  $(1/2, 3/4)$ . Thus, we can weaken the assumption (1.2)-(1.3) of theorem 1.1 appropriately.

## 5. Application.

As an example of application of theorem 1.1 we shall consider a model of magneto-micropolar fluid. In the incompressible case the governing system of equations for this model is [7]:

$$(5.1) \quad \operatorname{div} u = 0, \quad \operatorname{div} h = 0$$

$$(5.2) \quad \frac{\partial u}{\partial t} - (\nu + \nu_r) \Delta u + (u \cdot \nabla)u + \nabla(p + \frac{1}{2}h \cdot h) = 2\nu_r \operatorname{rot} \omega + rh \cdot \nabla h + f$$

$$(5.3) \quad j \frac{\partial \omega}{\partial t} - \alpha \Delta \omega - \beta \nabla \operatorname{div} \omega + j(u \cdot \nabla) \omega + 4\nu_r \omega = 2\nu_r \operatorname{rot} u + g$$

$$(5.4) \quad \frac{\partial h}{\partial t} - \gamma \Delta h + u \cdot \nabla h - h \cdot \nabla u = 0$$

where  $u$  is the velocity,  $p$  is the pressure,  $\omega$  is the microrotation (angular velocity of rotation of particles) and  $h$  is the magnetic field. Moreover,  $f$  and  $g$  are external fields and  $\nu, \nu_r, j, \alpha, \beta, \gamma$  are positive constants ( $\nu$  is the usual Newtonian viscosity,  $\nu_r$  is the microrotation viscosity). We assume that the density of the fluid is equal to one. We also set  $r = j = 1$  for simplicity. Let  $\Omega$  be a bounded set in  $R^3$  with smooth boundary.

By  $\mathcal{H}$  we denote the Hilbert space  $H \times L^2(\Omega)^3 \times H$  where  $H$  is the closure in the norm of  $L^2(\Omega)^3$  of the set of divergence free, smooth functions with compact support in  $\Omega$ . The norm in  $\mathcal{H}$  will be denoted by  $[\cdot]$ .

We introduce the following operators:

$$\mathcal{L}(U) = (-(\nu + \nu_r)\mathcal{P}\Delta u_1, -\alpha\Delta\omega_1 - (\alpha + \beta)\nabla\operatorname{div}\omega_1, -\gamma\mathcal{P}\Delta h_1) = (A_1u_1, A_2\omega_1, A_3h_1)$$

for  $U = (u_1, \omega_1, h_1) \in D(\mathcal{L})$ ,  $V = (u_2, \omega_2, h_2) \in D(\mathcal{L})$ ,

$$B(U, V) = (\mathcal{P}(u_1 \cdot \nabla u_2) - (\mathcal{P}h_1 \cdot \nabla h_2), u_1 \cdot \nabla \omega_2, (\mathcal{P}u_1 \cdot \nabla h_2) - (\mathcal{P}h_1 \cdot \nabla u_2))$$

$$R(U) = (-2\nu_r \operatorname{rot} \omega_1, -2\nu_r \operatorname{rot} u_1 + 4\nu_r \omega_1, 0)$$

The operator  $\mathcal{P}$  above is the orthogonal projection in  $L^2(\Omega)^3$  on the subspace  $H$ . In this notation the system of equations (5.1)-(5.4) takes the form (1.1) with  $F = (\mathcal{P}f, g, 0)$

We check the assumptions of theorem 1.1.

Operator  $\mathcal{L}$  is self-adjoint, positive with  $D(\mathcal{L}) = W^{2,2}(\Omega)^3 \times W^{2,2}(\Omega)^3 \times W^{2,2}(\Omega)^3 \cap (W_0^{1,2}(\Omega)^3 \times W_0^{1,2}(\Omega)^3 \times W_0^{1,2}(\Omega)^3) \cap \mathcal{H}$  where  $W^{2,2}(\Omega)$  and  $W_0^{1,2}(\Omega)$  are Sobolev spaces. The norms  $[\mathcal{L}U]$  and  $\|U\|_{W^{2,2}}$  are equivalent on  $D(\mathcal{L})$ . We denote  $X^\theta = D(\mathcal{L}^\theta)$  and  $Y_i^\theta = D(A_i^\theta)$ ,  $i = 1, 2, 3$ . Observe that  $Y_2^\theta = D(-\Delta^\theta)$ ,  $0 \leq \theta \leq 1$ .

Operator  $B$  satisfies (1.4), cf. [7]. From properties of the Laplace and the Stokes operators [3], [4], [5] we have, in particular,  $Y_i^{1/4} \subset L^3(\Omega)$ ,  $Y_i^{1/2} \subset L^6(\Omega)$ ,  $Y_i^{3/4} \subset W^{1,3}(\Omega)$ , with continuous imbeddings. Thus, for  $u \in Y_1^{1/4}$ ,  $v \in Y_i$ ,  $i = 1, 2, 3$ ,

$$|u \cdot \nabla v| \leq c|u|_{L^3} \cdot |\nabla v|_{L^6} \leq c_1|A_1^{1/4}u| \cdot |A_i v|$$

and, similarly, for  $u \in Y_1^{1/2}$ ,  $v \in Y_i^{3/4}$ ,  $i = 1, 2, 3$ ,

$$|u \cdot \nabla v| \leq c|u|_{L^6} \cdot |\nabla v|_{L^3} \leq c_1|A_1^{1/2}u| \cdot |A_i^{3/4}v|.$$

Finally, after simple calculations, we obtain

$$|B(U, V)| \leq c_1[\mathcal{L}^{1/4}U] \cdot [\mathcal{L}V]$$

and

$$|B(U, V)| \leq c_1[\mathcal{L}^{1/2}U] \cdot [\mathcal{L}^{3/4}V]$$

for  $U \in X^{1/4}$ ,  $V \in X^1$  and  $U \in X^{1/2}$ ,  $V \in X^{3/4}$ , respectively.

Moreover, operator  $R$  satisfies (1.5) and also (1.6) holds, cf. [6].

From theorem 1.1, cf. Remark 4.1, follows existence of  $\tau$ -periodic solution for the system of equations of magneto-micropolar fluids.

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