

# CRITICAL SUPERLINEAR AMBROSETTI-PRODI PROBLEMS

DJAIRO G. DE FIGUEIREDO AND YANG JIANFU  
IMECC-UNICAMP

## 1 INTRODUCTION

The main purpose of this work is to investigate the existence of multiple solutions of the critical superlinear problem

$$(1.1) \quad -\Delta u = \lambda u + u_+^{2^*-1} + f(x) \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega,$$

where  $2^* = \frac{2N}{N-2}$ ,  $N \geq 3$  is the critical Sobolev exponent, and  $\lambda > 0$  is a constant.  $u^+$  denotes the positive part of  $u$  :  $u^+(x) = \max\{u(x), 0\}$ .

This problem belongs to a class of problems which are known as the Ambrosetti-Prodi type. Due to the important role of the Ambrosetti-Prodi result [2] in subsequent research and for completeness we state it next. Let  $g : \mathbb{R} \rightarrow \mathbb{R}$  be a  $C^2$ -function such that  $g''(s) > 0$  for all  $s \in \mathbb{R}$  and

$$0 < \lim_{s \rightarrow -\infty} g'(s) < \lambda_1 < \lim_{s \rightarrow +\infty} g'(s) < \lambda_2,$$

where  $\lambda_1$  and  $\lambda_2$  are the first and second eigenvalues of  $(-\Delta, H_o^1(\Omega))$ . They consider the following boundary value problem

$$(1.2) \quad -\Delta u = g(u) + f(x) \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega,$$

where  $\Omega$  is a bounded domain in  $\mathbb{R}^N$  with a  $C^{2,\alpha}$  boundary  $\partial\Omega$ . Then, there is a  $C^1$  manifold  $M$  in  $C^{0,\alpha}(\bar{\Omega})$ , which splits the space into two open sets  $O_o$  and  $O_2$  with the following properties

- (i) if  $f \in O_o$ , problem (1.2) has no solution
- (ii) if  $f \in M$ , problem (1.2) has exactly one solution
- (iii) if  $f \in O_2$ , problem (1.2) has exactly two solutions.

A solution here means a function  $u \in C^{2,\alpha}(\bar{\Omega})$ .

After this work, several authors have extended this result in different directions. The literature on this problem is quite extensive; even risking the possibility of omitting some important work, we mention the following papers [1], [3], [4], [12], [17], [18] etc.

The above result shows the role that the location of the limits

$$(1.3) \quad g_- = \lim_{s \rightarrow -\infty} \frac{g(s)}{s}, \quad g_+ = \lim_{s \rightarrow +\infty} \frac{g(s)}{s}$$

with respect to the spectrum of  $(-\Delta, H_o^1(\Omega))$  plays in the question of existence of solutions for problem (1.2). Indeed, this contrasts with the well-known fact that if  $g_{\pm}$  are strictly between two consecutive eigenvalues, or both  $g_{\pm}$  are strictly less than  $\lambda_1$ , then problem (1.2) is solvable for all  $f$ . (We are assuming that  $f$  is locally Lipschitzian, and then solutions are in  $C^{2,\alpha}(\Omega) \cap C^0(\bar{\Omega})$ ). So the interesting cases are when the interval  $(g_-, g_+)$  contains eigenvalues. Problems with this feature are called problems of the Ambrosetti-Prodi type, or problems with jumping nonlinearities in a terminology introduced by Fučík, see [17]. These Ambrosetti-Prodi type of problems can be seen as a question of characterizing (or at least, describing part of) the range of a perturbation of a linear operator (say,  $-\Delta$ ) by some nonlinear operator (say  $Nu := -g(x, u)$ , which in our case is  $g(x, u) := \lambda u + u_+^{2^*-1}$ ). We can distinguish three different types of Ambrosetti-Prodi problems.

In type I, we have  $g_- < \lambda_1 < g_+$ , where  $g_-$  could be  $-\infty$ , and  $g_+$  could be  $+\infty$ . We write  $f = t\phi_1 + h$ , where  $t \in \mathbb{R}$ ,  $\phi_1$  is a first eigenfunction of  $(-\Delta, H_o^1(\Omega))$  with  $\phi_1 > 0$  and  $\int_{\Omega} \phi_1^2 dx = 1$ , and  $\int_{\Omega} h\phi_1 dx = 0$ . Then we can prove that in this case there is a  $t_o$  such that if  $t < t_o$ , problem (1.2) has at least one solution. Such a result holds under more general assumptions. Namely  $g$  can depend also on  $x$ , and the first limit in (1.3) can be replaced by limsup. Similarly the second limit can be replaced by liminf. See, for instance, the survey paper [16].

Type II is when  $g_-$  and  $g_+$  are finite, with the interval  $(g_-, g_+)$  containing eigenvalues. These problems are asymptotically linear. They have been extensively studied by Lazer-McKenna, see for instance [20]. In the treatment of this problem, via Topological and Variational Methods, it has appeared in an essential way the so-called Fučík spectrum, [17].

Type III is when  $g_{\pm}$  is between two consecutive eigenvalues and  $g_{\pm} < \lambda_1$ , these

are superlinear problems with a crossing of all but a finite number of eigenvalues. In this case one can prove that there is a  $t_o \in \mathbb{R}$  such that problem (1.2) with  $f = t\phi_1 + h$  has a negative solution for  $t > t_o$ . These problems have been treated in [25], and [15].

We remark that existence of a first solution for problems of type I and III does not require any growth at  $\pm\infty$ . So subcritical, critical or supercritical problems are treated. Observe that the reason is that: (i) in type I, one can find a subsolution and a supersolution, so a solution of problem (1.2) comes either by the Monotone Iteration Method if, for instance, the derivative of  $g$  is bounded, or by some Variational Methods after an appropriate truncation of the nonlinearity; (ii) in the case of type III we truncate the nonlinearity  $g$  for  $s > 0$ , getting a function  $\tilde{g}$  in such a way that  $g_-$  and  $\lim_{s \rightarrow +\infty} \frac{\tilde{g}(s)}{s}$  are between the same pair of consecutive eigenvalues.

The importance of the growth of  $g$  at infinite comes when one tries to get a second solution. The reason being that in order to have the functional associated to Equation (1.2)

$$I(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx - \int_{\Omega} G(u) dx - \int_{\Omega} f u dx$$

well defined in  $H_o^1(\Omega)$ , one has to require that

$$|g(s)| \leq C|s|^p + C,$$

where  $1 \leq p \leq 2^* - 1$ . The subcritical case  $p < 2^* - 1$  has been discussed by several authors mentioned before. Recently, Deng [13] considered problem (1.2) with a nonlinearity of the type  $g(u) = |u|^{2^*-1} + k(u)$ , where  $k$  is a lower perturbation of the expression with the critical exponent. This problem belongs to an Ambrosetti-Prodi problem of type I. In this case, the variational tool is the Mountain Pass Theorem.

Our problem stated in the beginning of this Introduction is of type I if  $\lambda < \lambda_1$  and of type III if  $\lambda > \lambda_1$ . In order to get a second solution, we have to recourse to a Linking Theorem. Both the geometry of functional associated to equation (1.2) and the determination of the levels where a  $(PS)$  condition fails are much more involved in type III than in type I. All along this paper we write the non-homogeneous term in the form  $f = t\phi_1 + h$ , where  $h \perp \phi_1$  in the  $L^2$ -sense. Let  $(0 <) \lambda_1 < \lambda_2 \leq \lambda_3 \leq \dots$  be eigenvalues of  $-\Delta$  subject to Dirichlet data, with corresponding eigenfunctions  $\phi_1, \phi_2, \phi_3, \dots$ . In Section 2, we prove the following result

**Theorem 1.1.** (I)(Existence of a negative solution)

(i) If  $0 < \lambda < \lambda_1$  and given  $h \in L^2$ , then there exists a  $t_o = t_o(h) < 0$  such that if  $t < t_o$ , Problem (1.1) has a negative solution  $u_t$ .

(ii) If  $\lambda > \lambda_1$ , and given  $h \in L^2$ , such that  $h \in \ker(-\Delta - \lambda)^\perp$  in the case that  $\lambda$  is an eigenvalue, then there exists  $t_o = t_o(h) > 0$  such that if  $t > t_o$ , Problem (1.1) has a negative solution  $u_t$ .

(II)(Existence of a second solution) If, in addition to either of the hypotheses above, one assumes that  $\lambda$  is not an eigenvalue of  $(-\Delta, H_o^1(\Omega))$  and the dimension  $N > 6$ , Problem (1.1) has a second solution.

Although the methods used here are essentially the same as for problems of Brézis-Nirenberg type namely

$$(1.4) \quad -\Delta u = |u|^{2^*-2}u + g(x, u) \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega,$$

where  $g(x, 0) = 0$  and  $g$  is some perturbation of lower order of the critical power, the technicalities have some new features. Indeed, for problem (1.4) the first solution is  $u \equiv 0$ , and from there one builds up the variational approach. In case of (1.2), the first solution  $u_t \neq 0$  and the translation of the functional to be centered at  $u_t$  introduces nonhomogeneities which are delicate to handle.

When one of the limits  $g_-$  or  $g_+$  is equal to an eigenvalue, we have a resonant problem. The solvability of (1.2) in this situation requires usually some additional conditions on  $g$ , like the Landesman-Lazer condition, see [20]. In Section 3 we discuss a case of resonance at  $\lambda = \lambda_1$ , where such a condition does not hold. Namely, the following result is proved.

**Theorem 1.2.** Suppose  $\lambda = \lambda_1$ . Then there is an  $\epsilon > 0$  such that if  $\|f\|_{L^2} < \epsilon$ , then (1.1) has a solution.

Finally in Section 4, we discuss local bifurcation at  $\lambda = \lambda_k, k > 1$ . Using the theory of bifurcation for variational problems as developed by Böhme [5] and Marino [21], we can handle eigenvalues of any algebraic multiplicity, and prove the next result.

**Theorem 1.3.** Let  $h \in \ker(-\Delta - \lambda_k)^\perp$  with  $k > 1$ . In the space  $\mathbb{R} \times H_o^1(\Omega)$ , let  $(\lambda, u_t(\lambda))$  for  $\lambda$  near  $\lambda_k$  be the line of negative solutions of (1.1) obtained in Theorem 1.1. Then  $(\lambda_k, u_t(\lambda_k))$  is a point of bifurcation

## 2 THE PROOF OF THEOREM 1.1

We write  $f(x) = t\phi_1(x) + h(x)$ , where  $\phi_1$  is the first eigenfunction of  $-\Delta$ ,  $\phi_1 \perp h$  in  $L^2$ -sense. We first prove that (1.1) has a negative solution  $u_t$ . Indeed, all negative solutions of (1.1) satisfies

$$(2.1) \quad -\Delta u = \lambda u + t\phi_1 + h.$$

If  $\lambda$  is an eigenvalue of  $(-\Delta, H_o^1(\Omega))$ , we suppose that  $h \in \ker(-\Delta - \lambda)^\perp$ . Then the problem

$$(2.2) \quad -\Delta u = \lambda u + h \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega$$

has a unique solution  $u_o$ . Consequently, the function  $w = u_t - u_o$ , where  $u_t$  is some solution of (2.1), is a solution of

$$(2.3) \quad -\Delta w = \lambda w + t\phi_1 \quad \text{in } \Omega, \quad w = 0 \quad \text{on } \partial\Omega.$$

Problem (2.3) has a unique solution  $w = \beta\phi_1$  where  $\beta = t/(\lambda_1 - \lambda)$ . Since we look for  $u_t \leq 0$ , it follows that: (i) for  $\lambda < \lambda_1$ , we obtain such  $u_t$  for  $t < 0$  and large, which comes from a negative  $\beta$ ; (ii) for  $\lambda > \lambda_1$ , we obtain such  $u_t$  for  $t > 0$  and large, which comes also from a negative  $\beta$ .

To find a second solution  $u$  of (1.1), we set  $u = v + u_t$ , and then  $v$  satisfies

$$(2.4) \quad -\Delta v = \lambda v + (v + u_t)_+^{2^*-1} \quad \text{in } \Omega, \quad v = 0 \quad \text{on } \partial\Omega.$$

So the second solution of (1.1) is obtained by finding a nontrivial solution  $v$  of (2.4).

Using variational methods we look for a critical point of the functional

$$I(v) = \frac{1}{2} \int_{\Omega} (|\nabla v|^2 - \lambda v^2) dx - \frac{1}{2^*} \int_{\Omega} (v + u_t)_+^{2^*} dx$$

defined in  $E = H_o^1(\Omega)$ . We use a Linking Theorem without Palais-Smale condition, see Theorem 4.3 in [22], or Theorem 5.1 in [14].

Suppose  $\lambda > 0$  is not an eigenvalue of  $(-\Delta, H_o^1(\Omega))$ . We assume  $\lambda \in (\lambda_k, \lambda_{k+1})$  from now on. The other case  $0 < \lambda < \lambda_1$  can be treated in a similar and simpler way, using Mountain Pass Theorem. Let us denote

$$E^- = \begin{cases} \emptyset & \text{if } \lambda \in (0, \lambda_1), \\ \{ \text{...} \} & \text{if } \lambda \in (\lambda_k, \lambda_{k+1}), \end{cases}$$

and

$$E^+ = (E^-)^\perp.$$

Let

$$S_\rho = \partial B_\rho \cap E^+$$

and

$$Q = [0, R] \oplus (\bar{B}_r \cap E^-), e \in E^+,$$

where  $\rho > 0, R > 0$  and  $r > 0$  will be determined later and in a way that

$$(2.5) \quad I|_{S_\rho} \geq \alpha > 0, \quad \rho < R,$$

$$(2.6) \quad I|_{\partial Q} < \alpha,$$

$$(2.7) \quad \max_Q I < \frac{1}{N} S^{\frac{N}{2}},$$

where  $S$  is the best Sobolev constant. Inequalities (2.5) - (2.6) will give the geometry of the functional  $I$  required by the Linking Theorem of Rabinowitz [24]. We will use it in the version without the assumption of Palais-Smale, see Theorem 4.3 in [22] or Theorem 5.1 in [14]. For that matter, condition (2.7) is used to prove that the solution obtained as a weak limit of a  $(PS)$ -sequence at the minimax level is not a trivial one.

**Lemma 2.1.** *There exist  $\rho_o > 0$  and a function  $\alpha > 0, \alpha : [0, \rho_o] \rightarrow \mathbb{R}$  such that*

$$I(v) \geq \alpha(\rho) \quad \text{for all } v \in S_\rho = \partial B_\rho \cap E^+.$$

*Explicitly, we have*

$$\rho_o = \left\{ S^{\frac{N}{N-2}} \left( 1 - \frac{\lambda}{\lambda_{k+1}} \right) \right\}^{\frac{N-2}{4}}, \quad \alpha(\rho) = \frac{1}{2} \left( 1 - \frac{\lambda}{\lambda_{k+1}} \right) \rho^2 - \frac{1}{2^*} S^{-\frac{N}{N-2}} \rho^{2^*}$$

*and the maximum value of  $\alpha(\rho)$*

$$\hat{\alpha} = \frac{1}{N} S^{\frac{N}{2}} \left( 1 - \frac{\lambda}{\lambda_{k+1}} \right)^{\frac{N}{2}}$$

*is assumed at*

$$\hat{\rho} = \left( 1 - \frac{\lambda}{\lambda_{k+1}} \right)^{\frac{N-2}{4}} S^{\frac{N}{4}}.$$

*Proof.* Using the fact that  $u_t < 0$  and the variational characterization of  $\lambda_{k+1}$  we get

$$I(v) \geq \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_{k+1}}\right) \int_{\Omega} |\nabla v|^2 dx - \frac{1}{2^*} \int_{\Omega} v_+^{2^*} dx.$$

By Sobolev imbedding we obtain

$$\begin{aligned} I(v) &\geq \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_{k+1}}\right) \int_{\Omega} |\nabla v|^2 dx - \frac{1}{2^*} S^{-\frac{N}{N-2}} \left(\int_{\Omega} |\nabla v|^2 dx\right)^{2^*/2} \\ &= \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_{k+1}}\right) \rho^2 - \frac{1}{2^*} S^{-\frac{N}{N-2}} \rho^{2^*}. \end{aligned}$$

The result follows by maximizing the function defined by the last equality.  $\square$

The best Sobolev constant  $S$  used above is defined by

$$(2.8) \quad S = \inf\{\|\nabla u\|_2^2 / \|u\|_{2^*}^2 : u \neq 0, u \in H^1(\mathbb{R}^N)\}$$

which is assumed by the functions

$$(2.9) \quad \psi_{\epsilon}(x) = \left(\frac{\epsilon \sqrt{N(N-2)}}{\epsilon^2 + |x|^2}\right)^{\frac{N-2}{2}}, \quad \epsilon > 0.$$

Let  $\xi \in C_o^1(\mathbb{R}^N)$  be a function such that  $\xi(x) = 1$  on  $B_{\frac{1}{2}}(0)$ ,  $\xi(x) = 0$  on  $\mathbb{R}^N \setminus B_1(0)$  and  $0 \leq \xi(x) \leq 1$  on  $\mathbb{R}^N$ . We may assume  $B_1(0) \subset \Omega$ . Let  $\phi_{\epsilon}(x) = \xi(x)\psi_{\epsilon}(x)$ , then we have following estimates.

**Lemma 2.2.** ([8])

$$(2.10) \quad \|\nabla \phi_{\epsilon}\|_2^2 = S^{\frac{N}{2}} + o(\epsilon^{N-2})$$

$$(2.11) \quad \|\phi_{\epsilon}\|_{2^*}^{2^*} = S^{\frac{N}{2}} + o(\epsilon^N)$$

$$(2.12) \quad \|\phi_{\epsilon}\|_2^2 = \begin{cases} K_1 \epsilon^2 + o(\epsilon^{N-2}) & \text{if } N \geq 5, \\ K_1 \epsilon^2 |\log \epsilon^2| + o(\epsilon^2) & \text{if } N = 4. \end{cases}$$

$$(2.13) \quad \|\phi_{\epsilon}\|_1 \leq K_2 \epsilon^{\frac{N+2}{2}}$$

and

$$(2.14) \quad \|\phi_{\epsilon}\|_{2^*-1}^{2^*-1} \leq K_3 \epsilon^{\frac{N-2}{2}},$$

where  $K_1 > 0, K_2 > 0$  and  $K_3 > 0$  are constants.

Denote by  $P_{\pm}$  the orthogonal projections of  $E$  onto  $E^{\pm}$  respectively. Using arguments as in [11], we can prove the following lemma:

**Lemma 2.3.**

$$(2.15) \quad \left| \int_{\Omega} [(P_+\phi_\epsilon)^{2^*} - \phi_\epsilon^{2^*}] dx \right| \leq C\epsilon^{N-2},$$

$$(2.16) \quad \left| \int_{\Omega} (|\nabla\phi_\epsilon|^2 - |\nabla(P_+\phi_\epsilon)|^2) dx \right| \leq C\epsilon^{N-2},$$

$$(2.17) \quad \|P_+\phi_\epsilon\|_{2^*-1}^{2^*-1} \leq C\epsilon^{\frac{N-2}{2}}$$

$$(2.18) \quad \|P_+\phi_\epsilon\|_1 \leq C\epsilon^{\frac{N+2}{2}}$$

and

$$(2.19) \quad \|P_-\phi_\epsilon\|_\infty \leq C\epsilon^{\frac{N-2}{2}}.$$

Define for any fixed  $K > 0$  the set  $\Omega_{\epsilon,K} = \{x \in \Omega : P_+\phi_\epsilon(x) > K\}$ . By (2.19) we know that

$$P_+\phi_\epsilon(0) = \phi_\epsilon - P_-\phi_\epsilon(0) \geq C\epsilon^{-\frac{N-2}{2}} - \|P_-\phi_\epsilon\|_\infty \geq C\epsilon^{-\frac{N-2}{2}}$$

which implies  $P_+\phi_\epsilon(0) \rightarrow +\infty$  as  $\epsilon \rightarrow 0$ . By the continuity of  $P_+\phi_\epsilon$ , there exists  $\delta > 0$  such that  $B_\delta(0) \subset \Omega_{\epsilon,K}$ . Therefore we have a result as follows.

**Lemma 2.4.**

$$(2.20) \quad \int_{\Omega_{\epsilon,K}} |P_+\phi_\epsilon|^{2^*} dx = \int_{\Omega} |\phi_\epsilon|^{2^*} dx + O(\epsilon^{N-2}),$$

$$(2.21) \quad \int_{\Omega_{\epsilon,K}} |P_+\phi_\epsilon|^{2^*-1} dx = \int_{\Omega} \phi_\epsilon^{2^*-1} dx + O(\epsilon^{\frac{N+2}{2}}),$$

and

$$(2.22) \quad \int_{\Omega_{\epsilon,K}} |P_+\phi_\epsilon| dx = \int_{\Omega} \phi_\epsilon dx + O(\epsilon^N).$$



**Lemma 2.5.** *Let  $u, v \in L^p(\Omega)$  with  $2 \leq p \leq 2^*$ . If  $\omega \subset \Omega$  and  $u + v > 0$  on  $\omega$ , then*

$$(2.23) \quad \left| \int_{\omega} (u + v)^p dx - \int_{\omega} |u|^p dx - \int_{\omega} |v|^p dx \right| \leq C \int_{\omega} (|u|^{p-1}|v| + |u||v|^{p-1}) dx,$$

where  $C$  depends only on  $p$ .

*Proof.* By the Fundamental Theorem of Calculus the left side of (2.23) is equal to

$$\left| p \int_0^1 d\tau \int_{\omega} [ |v + \tau u|^{p-2}(v + \tau u) - |\tau u|^{p-2}\tau u ] u dx \right|,$$

which by its turn is equal to, using the mean value theorem

$$p(p-1) \left| \int_0^1 d\tau \int_{\omega} |\tau u + v\theta(x)|^{p-2} uv dx \right|, \quad 0 < \theta(x) < 1.$$

This last expression can be estimated by

$$C \int_0^1 d\tau \int_{\omega} (\tau^{p-2}|u|^{p-1}|v| + |u||v|^{p-1}) dx \leq C \int_{\omega} (|u|^{p-2}|v| + |u||v|^{p-1}) dx.$$

□

**Lemma 2.6.** *let  $A, B, C$  and  $\alpha$  be positive numbers. Consider the function*

$$\Phi_{\epsilon}(s) = \frac{1}{2}s^2 A - \frac{1}{2^*}s^{2^*} B + s^{2^*} \epsilon^{\alpha} C, \quad s > 0.$$

*Then*

$$s_{\epsilon} = \left( \frac{A}{B - 2^* \epsilon^{\alpha} C} \right)^{\frac{1}{2^*-2}}$$

*is the point where  $\Phi_{\epsilon}$  achieves its maximum. Write  $s_{\epsilon} = (1 + t_{\epsilon})s_0$ , where  $s_0 = \left(\frac{A}{B}\right)^{\frac{1}{2^*-2}}$  is the point at which  $\Phi_0$  achieves its maximum. Then*

$$t_{\epsilon} = O(\epsilon^{\alpha})$$

*and*

$$\Phi_{\epsilon}(s) \leq \Phi_{\epsilon}(s_{\epsilon}) = \frac{1}{2} \left( \frac{A^N}{B^{N-2}} \right)^{\frac{1}{2}} + O(\epsilon^{\alpha}).$$

*Proof.* It is clear that  $\Phi_{\epsilon}$  achieves its maximum at  $s_{\epsilon}$  and  $s_{\epsilon}$  satisfies

$$(2.24) \quad \frac{1}{2} A - \frac{2^*-1}{2^*} B + 2^* C \epsilon^{\alpha} s_{\epsilon}^{2^*-1} = 0$$

This implies

$$(2.25) \quad s_\epsilon \geq s_o.$$

Writing  $s_\epsilon = (1 + t_\epsilon)s_o$ , we derive from (2.24) that

$$(2.26) \quad s_\epsilon \rightarrow s_o, \quad t_\epsilon \rightarrow 0 \quad \text{as} \quad \epsilon \rightarrow 0$$

and

$$(2.27) \quad (1 + t_\epsilon)s_o A - (1 + t_\epsilon)^{2^*-1} s_o^{2^*-1} B + 2^* C \epsilon^\alpha (1 + t_\epsilon)^{2^*-1} s_o^{2^*-1} = 0.$$

That is

$$\left(\frac{A^{2^*-1}}{B}\right)^{\frac{1}{2^*-2}} [(1 + t_\epsilon) - (1 + t_\epsilon)^{2^*-1}] + 2^* C \epsilon^\alpha (1 + t_\epsilon)^{2^*-1} s_o^{2^*-1} = 0.$$

Expanding for  $t_\epsilon$  we obtain

$$(2.28) \quad \left[\frac{4}{N-2} t_\epsilon + o(t_\epsilon)\right] \left(\frac{A^{2^*-1}}{B}\right)^{\frac{1}{2^*-2}} = 2^* C \epsilon^\alpha (1 + t_\epsilon)^{2^*-1} s_o^{2^*-1}.$$

Hence

$$(2.29) \quad t_\epsilon = O(\epsilon^\alpha).$$

□

Our aim is to choose  $Q$  and  $\rho$  such that (2.5), (2.6) and (2.7) hold. So choose  $e$  as a function of  $\epsilon$ :  $e_\epsilon = P_+ \phi_\epsilon$ .

**Lemma 2.7.** *There exist  $r_o > 0, R_o > 0$ , and  $\epsilon_o > 0$  such that for  $r \geq r_o, R \geq R_o$  and  $0 < \epsilon \leq \epsilon_o$  we have*

$$I|_{\partial Q} < \alpha,$$

where  $\alpha > 0$  is determined in Lemma 2.1.

*Proof.* We may write  $\partial Q = \Gamma_1 \cap \Gamma_2 \cap \Gamma_3$  with

$$\Gamma_1 = \bar{B}_r \cap E^-,$$

$$\Gamma_2 = \{v \in E : v = w + s e_\epsilon, w \in E^-, \|w\| = r, 0 \leq s \leq R\},$$

$$\Gamma_3 = \{v \in E : v = w + s e_\epsilon, w \in E^- \cap B(0, r)\}$$

We will show that on each  $\Gamma_i$ , we have  $I|_{\Gamma_i} < \alpha, i = 1, 2, 3$ .

For any  $v \in E^-$  we have

$$(2.30) \quad \int_{\Omega} |\nabla v|^2 dx \leq \lambda_k \int_{\Omega} v^2 dx.$$

So for  $v \in \Gamma_1$

$$I(v) = \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_k}\right) \int_{\Omega} |\nabla v|^2 dx - \frac{1}{2^*} \int_{\Omega} (v + u_t)_+^{2^*} dx \leq 0.$$

For  $v \in \Gamma_2$ , we distinguish two cases.

Define  $\delta^2 = \sup_{0 < \epsilon \leq 1} \int_{\Omega} |\nabla e_{\epsilon}|^2 dx$ . If  $0 \leq s \leq s_o := \frac{\sqrt{2\hat{\alpha}}}{\delta}$ , then

$$\begin{aligned} I(v) &\leq \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_k}\right) r^2 + \frac{1}{2} s^2 \int_{\Omega} |\nabla e_{\epsilon}|^2 dx - \frac{1}{2^*} \int_{\Omega} (w + s e_{\epsilon} + u_t)_+^{2^*} dx \\ &\leq \frac{1}{2} s^2 \delta^2 < \hat{\alpha}. \end{aligned}$$

If  $s \geq s_o = \frac{\sqrt{2\hat{\alpha}}}{\delta}$ , denote

$$K = \sup \left\{ \left\| \frac{w + u_t}{s} \right\|_{L^{\infty}} : s_o \leq s \leq R, \|w\|_E = r, w \in E^- \right\}.$$

$K$  is independent of  $R$ . Since  $P_+ \phi_{\epsilon}(0) \rightarrow \infty$  as  $\epsilon \rightarrow 0$ , there exists  $\epsilon'_o > 0$  such that for all  $\epsilon, 0 < \epsilon < \epsilon'_o$  and  $s \geq s_o$

$$\Omega_{\epsilon} = \{x \in \Omega : e_{\epsilon}(x) > K\} \neq \emptyset.$$

Whence by Lemma 2.5

$$\begin{aligned} (2.31) \quad &\int_{\Omega} \left(e_{\epsilon} + \frac{w + u_t}{s}\right)_+^{2^*} dx = \int_{\Omega_{\epsilon}} \left(e_{\epsilon} + \frac{w + u_t}{s}\right)_+^{2^*} dx \\ &\geq \int_{\Omega_{\epsilon}} |e_{\epsilon}|^{2^*} dx + \int_{\Omega_{\epsilon}} \left|\frac{w + u_t}{s}\right|^{2^*} dx - C \int_{\Omega_{\epsilon}} \left(|e_{\epsilon}|^{2^*-1} \left|\frac{w + u_t}{s}\right| + |e_{\epsilon}| \left|\frac{w + u_t}{s}\right|^{2^*-1}\right) dx \\ &\geq \int_{\Omega_{\epsilon}} |e_{\epsilon}|^{2^*} dx + \int_{\Omega_{\epsilon}} \left|\frac{w + u_t}{s}\right|^{2^*} dx - C \left(\|e_{\epsilon}\|_{L^{2^*-1}(\Omega_{\epsilon})}^{2^*-1} + \|e_{\epsilon}\|_{L^1(\Omega_{\epsilon})}\right). \end{aligned}$$

By Lemmas 2.2, 2.3 and 2.4 and (2.31) we obtain

$$(2.32) \quad I(v) \leq \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_k}\right) r^2 + \frac{1}{2} s^2 S^{\frac{N}{2}} - \frac{s^{2^*}}{2^*} S^{\frac{N}{2}} + C s^{2^*} \epsilon^{\frac{N-2}{2}} := \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_1}\right) r^2 + \Phi_{\epsilon}(s).$$

Applying Lemma 2.6 to  $\Phi_{\epsilon}(s)$ , we obtain

$$I(v) < \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_1}\right) r^2 + \frac{1}{2} S^{\frac{N}{2}} + O(\epsilon^{\frac{N-2}{2}}).$$

We may choose  $r > 0$  such that  $I(v) < 0$ . This determines  $r_o$ .

For  $v \in \Gamma_3$  we have  $v = w + Re_\epsilon$ ,  $w \in E^- \cap B_r(0)$  and

$$I(v) \leq \frac{1}{2} \left(1 - \frac{\lambda}{\lambda_k}\right) \|w\|_E^2 + \frac{1}{2} R^2 \int_{\Omega} |\nabla e_\epsilon|^2 dx - \frac{1}{2^*} R^{2^*} \int_{\Omega} \left(e_\epsilon + \frac{u_t + w}{R}\right)_+^{2^*} dx.$$

By the boundedness of  $w$  and  $u_t$ , there exists  $K > 0$  such that

$$\|w + u_t\|_{L^\infty(\Omega)} \leq K.$$

Again since  $P_+ \phi_\epsilon(0) \rightarrow +\infty$  as  $\epsilon \rightarrow 0$ , there exists  $\epsilon_o > 0$  (take  $\epsilon_o < \epsilon'_o$ ) such that if  $0 < \epsilon < \epsilon_o$ ,  $P_+ \phi_\epsilon(0) > 2K$ . Given  $\epsilon > 0$ ,  $0 < \epsilon < \epsilon_o$ , there exist  $R_o = R_o(\epsilon)$ ,  $\eta = \eta(\epsilon)$  such that

$$|\{x \in \Omega : e_\epsilon + \frac{w + u_t}{R} > 1\}| \geq \eta > 0$$

for all  $R > R_o$ . Hence we find  $\epsilon_o, R_o > 0$  such that if  $\epsilon < \epsilon_o$ , and  $R > R_o$ , we have  $I(v) \leq 0$  for  $v \in \Gamma_3$ .

□

**Lemma 2.8.**

$$(2.33) \quad \max_Q I < \frac{1}{N} S^{\frac{N}{2}}.$$

*Proof.* Let us fix  $\epsilon < \epsilon_o$ , so that the geometry of the Linking Theorem holds. For  $w + se_\epsilon \in Q$ , we have

$$(2.34) \quad I(w + se_\epsilon) = \frac{1}{2} \int_{\Omega} (|\nabla w|^2 - \lambda w^2) dx + \frac{1}{2} s^2 \int_{\Omega} (|\nabla e_\epsilon|^2 - \lambda e_\epsilon^2) dx - \frac{1}{2^*} \int_{\Omega} (w + se_\epsilon + u_t)_+^{2^*} dx. \quad \blacksquare$$

With the same notations and arguments as in the proof of Lemma 2.7, if  $s < s_o$  we have

$$(2.35) \quad I(w + se_\epsilon) \leq \frac{1}{2} s^2 \int_{\Omega} |\nabla e_\epsilon|^2 dx = \frac{1}{2} s^2 \delta^2 < \frac{1}{N} S^{\frac{N}{2}}.$$

If  $s \geq s_o$ , using (2.31) we deduce as (2.32) that

$$(2.36) \quad I(w + se_\epsilon) \leq \frac{1}{2} s^2 \int_{\Omega} (|\nabla e_\epsilon|^2 - \lambda e_\epsilon^2) dx - \frac{1}{2^*} s^{2^*} \int_{\Omega} e_\epsilon^{2^*} dx + C s^{2^*} \epsilon^{\frac{N-2}{2}} := \Phi_\epsilon(s).$$

An application of Lemma 2.6 to  $\Phi_\epsilon(s)$  yields

$$I(w + se_\epsilon) = \frac{1}{N} \left[ \int_{\Omega} (|\nabla e_\epsilon|^2 - \lambda e_\epsilon^2) dx \right]^{\frac{N}{2}} \left( \int_{\Omega} e_\epsilon^{2^*} dx \right)^{-\frac{N-2}{2}} + O(\epsilon^{\frac{N-2}{2}}).$$

Using the estimates in Lemmas 2.2, 2.3 and 2.4 on  $e_\epsilon$  we get

$$I(w + se_\epsilon) \leq \frac{1}{N} S^{\frac{N}{2}} - \frac{1}{2} \lambda \begin{cases} O(\epsilon^2), & \text{if } N \geq 5 \\ O(\epsilon^2 |\log \epsilon^2|), & \text{if } N = 4 \end{cases} + O(\epsilon^{\frac{N-2}{2}}).$$

If  $N > 6$ , i.e.  $2 < \frac{N-2}{2}$ , the result follows by choosing  $\epsilon > 0$  sufficiently small.

□

*Proof of Theorem 1.1.* It remains to prove the existence of a second solution of (1.1), i.e. a nontrivial solution of (2.4). Using the Linking Theorem, Lemmas 2.1 and 2.7, there exists  $\{v_n\} \subset E$  such that

$$(2.37) \quad I(v_n) = \frac{1}{2} \int_{\Omega} (|\nabla v_n|^2 - \lambda v_n^2) dx - \frac{1}{2^*} \int_{\Omega} (v_n + u_t)_+^{2^*} dx = c + o(1),$$

$$(2.38) \quad \langle I'(v_n), \phi \rangle = \int_{\Omega} (\nabla v_n \nabla \phi - \lambda v_n \phi) dx - \int_{\Omega} (v_n + u_t)_+^{2^*-1} \phi dx = o(1) \|\phi\|_E,$$

where  $c$  is the minimax level in the Linking Theorem with  $e_\epsilon = P_+ \phi_\epsilon$ , and  $\epsilon < \epsilon_0$  sufficiently small in order to have the validity of Lemmas 2.7 and 2.8, and  $Q$  as above.

First we prove that  $\{v_n\}$  is bounded in  $E$ . It follows from (2.37) and (2.38)

$$\frac{1}{N} \int_{\Omega} (v_n + u_t)_+^{2^*} dx - \frac{1}{2} \int_{\Omega} (v_n + u_t)_+^{2^*-1} u_t dx \leq c + \epsilon_n \|v_n\|_E + o(1)$$

where  $\epsilon_n \rightarrow \infty$  as  $n \rightarrow \infty$ . It implies

$$(2.39) \quad \int_{\Omega} (v_n + u_t)_+^{2^*} dx \leq c + \epsilon_n \|v_n\|_E + o(1)$$

since  $u_t \leq 0$ . Writing  $v_n = v_n^+ + v_n^-$ , with  $v_n^\pm \in E^\pm$  we get from (2.37) - (2.38), using Hölder and Young inequalities that

$$\begin{aligned} & \left(1 - \frac{\lambda}{\lambda_{k+1}}\right) \|v_n^+\|_E^2 \\ & \leq \int_{\Omega} (v_n + u_t)_+^{2^*-1} v_n^+ dx + \epsilon \|v_n^+\| \\ & \leq \epsilon \left( \int_{\Omega} |v_n^+|^{2^*} dx \right)^{\frac{2}{2^*}} + C_\epsilon \left( \int_{\Omega} (v_n + u_t)_+^{2^*} dx \right)^{\frac{2(2^*-1)}{2^*}} + \epsilon_n \|v_n^+\|_E \\ & \leq \epsilon \left( \int_{\Omega} |v_n^+|^{2^*} dx \right)^{\frac{2}{2^*}} + C_\epsilon + \epsilon_n \left( \|v_n\|_E^{\frac{N+2}{N}} + \|v_n^+\|_E \right). \end{aligned}$$

So we obtain

$$\|v_n^+\|_E^2 \leq C + \epsilon \left( \|v_n\|_E^{\frac{N+2}{N}} + \|v_n^+\|_E \right)$$

In the same way, we have

$$\|v_n^-\|_E^2 \leq C + \epsilon_n (\|v_n\|_E^{\frac{N+2}{N}} + \|v_n^-\|_E).$$

Consequently,

$$\|v_n\|_E \leq C.$$

Hence we may assume

(2.40)

$$v_n \rightarrow v \text{ weakly in } H_o^1(\Omega), \quad v_n \rightarrow v \text{ in } L^q(\Omega) \quad 2 \leq q < 2^*, \quad v_n \rightarrow v \text{ a.e. in } \Omega,$$

as  $n \rightarrow \infty$ . It follows that  $v$  is a weak solution of

$$(2.41) \quad -\Delta v = \lambda v + (v + u_t)_+^{2^*-1}$$

which implies

$$(2.42) \quad \int_{\Omega} (|\nabla v|^2 - \lambda v^2) dx - \int_{\Omega} (v + u_t)_+^{2^*} dx + \int_{\Omega} (v + u_t)_+^{2^*-1} u_t dx = 0.$$

By Brézis - Lieb Lemma [7]

$$(2.43) \quad \int_{\Omega} (v_n + u_t)_+^{2^*} dx = \int_{\Omega} (v_n - v)_+^{2^*} dx + \int_{\Omega} (v + u_t)_+^{2^*} dx + o(1).$$

Hence, using (2.43)

$$(2.44) \quad I(v_n) = I(v) + \frac{1}{2} \int_{\Omega} |\nabla(v_n - v)|^2 dx - \frac{1}{2^*} \int_{\Omega} (v - v_n)_+^{2^*} dx + o(1),$$

and similarly by (2.41)

(2.45)

$$\langle I'(v_n), v_n \rangle = \int_{\Omega} |\nabla(v_n - v)|^2 dx - \int_{\Omega} (v_n - v)_+^{2^*} dx - \int_{\Omega} (v_n - v)_+^{2^*-1} u_t dx + o(1).$$

Since  $\int_{\Omega} (v_n - v)_+^{2^*-1} u_t dx \rightarrow 0$  as  $n \rightarrow \infty$ , it yields

$$(2.46) \quad \int_{\Omega} |\nabla(v_n - v)|^2 dx = \int_{\Omega} (v_n - v)_+^{2^*} dx + o(1).$$

Let  $w_n = v_n - v$  and

$$(2.47) \quad \lim_{n \rightarrow \infty} \int_{\Omega} |\nabla w_n|^2 = k \geq 0.$$

If  $k = 0$ , then  $v_n \rightarrow v$  strongly in  $H_o^1(\Omega)$  as  $n \rightarrow \infty$ , then  $\alpha \leq c = I(v)$ ,  $v$  is a nontrivial solution of (2.4).

If  $k > 0$ , we claim that  $v \neq 0$ . Indeed, using (2.46) and the Sobolev inequality we obtain

$$(2.48) \quad \|w_n\|_{H_o^1}^2 \geq S \left( \int_{\Omega} |w_n|^{2^*} dx \right)^{2/2^*} \geq S \left( \int_{\Omega} (w_n)_+^{2^*} dx \right)^{2/2^*} \geq S \left[ \int_{\Omega} |\nabla w_n|^2 dx + o(1) \right]^{2/2^*}$$

which gives

$$(2.49) \quad k \geq S k^{\frac{N-2}{N}} \quad \text{i.e.} \quad k \geq S^{\frac{N}{2}}.$$

From (2.44), (2.46) and (2.49), if  $v \equiv 0$  we have

$$c + o(1) = \frac{k}{N} \geq \frac{1}{N} S^{\frac{N}{2}}.$$

It contradicts to the statement of Lemma 2.8. Therefore  $v \neq 0$ . By (2.41) we know  $v$  is not negative.

□

### 3. EXISTENCE OF SOLUTIONS FOR THE CASE $\lambda = \lambda_1$

We consider

$$(3.1) \quad \begin{cases} -\Delta u = \lambda_1 u + u_+^{2^*-1} + f & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega. \end{cases}$$

A necessary condition for the solvability of (3.1) is given by

$$(3.2) \quad \int_{\Omega} f \phi_1 dx < 0,$$

where  $\phi_1$  is the first eigenfunction of  $-\Delta$ . Although one expects that (3.2) would be a sufficient condition for the solvability of (3.1), we have not been able to prove it. Indeed, we require in addition that  $f$  has small  $L^2$ -norm. Let  $E^- = \text{span}\{\phi_1\}$  and  $E^+ = (E^-)^\perp$ . For any  $u \in E$  there are  $t \in \mathbb{R}$  and  $v \in E$  such that  $u = t\phi_1 + v$ . The functional  $I : E \rightarrow \mathbb{R}$  associated with equation (3.1) can be written as

$$I(u) = \frac{1}{2} \int_{\Omega} [|\nabla v|^2 - \lambda_1 v^2] dx - \frac{1}{2^*} \int_{\Omega} (v + t\phi_1)_+^{2^*} dx - \int_{\Omega} f(v + t\phi_1) dx,$$

where  $\phi_1 = \phi_1 + t\phi_1 = \phi_1 + t\phi_1$ .

**Lemma 3.1.** *For any given  $v \in E^+$ , the functional  $I$  is bounded above in  $E^-$ .*

*Proof.* Given  $v \in E^+$ , let us define the real-valued function

$$(3.3) \quad g(t) = I(v + t\phi_1).$$

For  $t < 0$  we have

$$g(t) \leq \frac{1}{2} \int_{\Omega} [|\nabla v|^2 - \lambda_1 v^2] dx + \|f\|_{L^2} \|v\|_{L^2}.$$

For  $t > 0$ , we claim

$$(3.4) \quad \lim_{t \rightarrow +\infty} \left\{ \frac{1}{2^*} \int_{\Omega} (v + t\phi_1)_+^{2^*} dx + \int_{\Omega} f(v + t\phi_1) dx \right\} = +\infty$$

which completes the proof, since  $g$  is continuous.

To prove (3.4) we proceed as follows: let  $a = \max\{\phi_1(x) : x \in \Omega\}$ . Choose  $\Omega_o \subset\subset \Omega$  such that  $\phi_1(x) > \frac{a}{2}$  for  $x \in \Omega_o$ . By Lusin's theorem, given any  $\delta > 0$  (choose  $\delta = \frac{1}{2}|\Omega_o|$ ) there exists a continuous function  $h(x)$  in  $\Omega_o$  such that

$$\text{meas}\{x : h(x) \neq v(x)\} < \delta.$$

So the set  $G = \{x : h(x) = v(x)\}$  has measure greater than  $|\Omega_o| - \delta$ . Let  $M = \sup\{|v(x)| : x \in G\}$ . Then, for  $x \in G$  we have

$$\phi_1(x) + \frac{v(x)}{t} \geq \frac{a}{2} - \frac{M}{t} \geq \frac{a}{4}$$

if  $t \geq t_o := \frac{4M}{a}$ . So there is  $\eta > 0$  such that

$$\int_{\Omega} \left(\phi_1 + \frac{v}{t}\right)_+^{2^*} dx \geq \eta \quad \text{for } t \geq t_o.$$

Then the first term in (3.4) is larger than  $Ct^{2^*}$  which proves the claim.

□

Next we claim that, for each  $v \in E^+$ , there is a unique  $t(v)$  such that

$$(3.5) \quad g(t(v)) = \max\{g(t) : t \in \mathbb{R}\}.$$

At a point  $t_o$  of maximum we have  $g'(t_o) = 0$ , i.e.

$$(3.6) \quad g'(t_o) = - \int_{\Omega} (t_o \phi_1 + v)_+^{2^*-1} \phi_1 dx - \int_{\Omega} f \phi_1 dx = 0.$$



If  $t \geq t_o$ ,  $g'(t) \leq g'(t_o)$  and if  $t \leq t_o$ ,  $g'(t) \geq g'(t_o)$ , hence

$$g(t_o) = \max_{\{t \in \mathbb{R}\}} g(t), \text{ i.e. } I(t\phi_1 + v) \leq I(t_o\phi_1 + v).$$

The second derivative of  $g$  is given by

$$g''(t) = - \int_{\Omega} (t\phi_1 + v)_+^{2^*-2} \phi_1^2 dx \leq 0$$

which says that  $g$  is concave. So the set of maxima is a closed interval, and we show it is a single point. At a point  $t_o$  of maximum  $g''(t_o)$  can not be 0. Indeed, if this were the case, then  $(t_o\phi_1 + v)_+ = 0$ , which would imply by (3.6) that  $\int_{\Omega} f\phi_1 dx = 0$ , a contradiction. So  $g$  is strictly concave at  $t_o$ . This also proves, as a consequence of the Implicit Function Theorem that the mapping

$$v \in E^+ \rightarrow t(v) \in \mathbb{R}$$

is continuous and differentiable. Therefore

$$(3.7) \quad I(t\phi_1 + v) \leq I(t(v)\phi_1 + v) \quad \text{if } t \neq t(v)$$

and from (3.6) we have

$$(3.8) \quad \int_{\Omega} (v + t(v)\phi_1)_+^{2^*-1} \phi_1 + \int_{\Omega} f\phi_1 = 0, \quad \forall v \in E^+.$$

The relation (3.6) for  $v = 0$  gives

$$(3.9) \quad \int_{\Omega} (t(0)\phi_1)_+^{2^*-1} \phi_1 dx = - \int_{\Omega} f\phi_1 dx$$

and the function  $g(t)$  in this case is

$$(3.10) \quad -\frac{1}{2^*} \int_{\Omega} (t\phi_1)_+^{2^*} dx - t \int_{\Omega} f\phi_1 dx$$

which shows that  $t(0)$  has to be greater than 0. So the relation (3.9) can be written as

$$(3.11) \quad t(0)^{2^*-1} \int_{\Omega} \phi_1^{2^*} dx = - \int_{\Omega} f\phi_1 dx.$$

Let us introduce the notations

$$(3.12) \quad A = - \int_{\Omega} f\phi_1 dx, \quad B = \int_{\Omega} \phi_1^{2^*} dx.$$

Our next step is to show that the functional  $F : E^+ \rightarrow \mathbb{R}$  given by

$$F(v) = I(v + t(v)\phi_1)$$

has a minimum in the interior of certain ball  $B_\rho$  centered at the origin.

It is easy to see that

$$(3.13) \quad F(0) = \frac{N+2}{2N} \frac{A^{2N/(N+2)}}{B^{(N-2)/(N+2)}},$$

and next we estimate  $F(v)$ :

$$(3.14) \quad F(v) = \frac{1}{2} \int_{\Omega} [|\nabla v|^2 - \lambda_1 v^2] dx - \frac{1}{2^*} \int_{\Omega} (v + t(v)\phi_1)_+^{2^*} dx - \int_{\Omega} f(v + t(v)\phi_1) dx.$$

Let

$$(3.15) \quad M_1 = \frac{1}{N+1} \lambda_2^{-\frac{N}{4}} S^{\frac{N}{4}} \left(\frac{N}{N+2}\right)^{\frac{N-2}{4}} (\lambda_2 - \lambda_1)^{\frac{N+2}{4}},$$

$$(3.16) \quad M_2 = \min\left\{\left(\frac{2}{N+2}\right)^{\frac{N+2}{2N}} S^{\frac{N+2}{4}}, \left(\frac{2}{N+2}\right)^{\frac{N+2}{2N}} \|\phi_1\|_{2^*} \left[\frac{N}{N+2} \left(1 - \frac{\lambda_1}{\lambda_2}\right) S\right]^{\frac{N+2}{4}}\right\}.$$

In addition to (3.2), we suppose that  $f$  satisfies

$$(3.17) \quad \|f\|_2 \leq M_1 \quad \text{and} \quad - \int_{\Omega} f \phi_1 dx < M_2.$$

**Lemma 3.2.** *Suppose (3.2) and (3.17), there is an  $\alpha > 0$  such that*

$$(3.18) \quad F(v) \geq \alpha > F(0)$$

provide that  $\|v\|_E = \rho_o$  with  $\rho_o = \left[\frac{N}{N+2} \left(1 - \frac{\lambda_1}{\lambda_2}\right)\right]^{\frac{N-2}{4}} S^{\frac{N}{4}}$ .

*Proof.* It follows from (3.6) and the inequality

$$\int_{\Omega} |\nabla v|^2 dx \geq \lambda_2 \int_{\Omega} v^2 dx, \quad \forall v \in E^+$$

that

$$(3.19) \quad \begin{aligned} F(v) &\geq I(v) = \frac{1}{2} \int_{\Omega} [|\nabla v|^2 - \lambda_1 v^2] dx - \frac{1}{2^*} \int_{\Omega} v_+^{2^*} dx - \int_{\Omega} f v dx \\ &\geq \frac{1}{2} \left(1 - \frac{\lambda_1}{\lambda_2}\right) \int_{\Omega} |\nabla v|^2 dx - \frac{1}{2^*} \int_{\Omega} |v|^{2^*} dx - \|f\|_2 \|v\|_2. \end{aligned}$$

Using Sobolev inequality and (3.19) we obtain

$$(3.20) \quad F(v) \geq \frac{1}{2} \left(1 - \frac{\lambda_1}{\lambda_2}\right) \rho - \frac{1}{2^*} S^{-\frac{N}{N-2}} \rho^{2^*} - \|f\|_2 \lambda_2^{-\frac{1}{2}} \rho,$$

where  $S$  is the best Sobolev constant and  $\rho = \left(\int_{\Omega} |\nabla v|^2 dx\right)^{\frac{1}{2}}$ . Consider the real function

$$k(\rho) = \frac{1}{2} a \rho^2 - \frac{1}{2^*} b \rho^{2^*} - c \rho := \rho j(\rho).$$

The maximum point  $\rho$  of  $j(\rho)$  on  $\mathbb{R}_+$  satisfies

$$j'(\rho_o) = \frac{1}{2} a - \frac{2^* - 1}{2^*} b \rho_o^{2^* - 2} = 0.$$

Then we have

$$\rho_o = \left[ \frac{N}{N+2} \left(1 - \frac{\lambda_1}{\lambda_2}\right) \right]^{\frac{N-2}{4}} S^{\frac{N}{4}}$$

and

$$(3.21) \quad k(\rho_o) = \rho_o \left[ \frac{2}{N+2} \left( \frac{N}{(N+2)b} \right)^{\frac{N-2}{4}} a^{\frac{N+2}{4}} - c \right].$$

With  $a = 1 - \frac{\lambda_1}{\lambda_2}$ ,  $b = S^{-\frac{N}{N-2}}$  and  $c = \|f\|_2 \lambda_2^{-\frac{1}{2}}$  in (3.21) and by the assumption (3.17) we obtain

$$(3.22) \quad F(v) \geq \frac{\rho_o}{N+2} \left[ \frac{N}{(N+2)b} \right]^{\frac{N-2}{4}} a^{\frac{N+2}{4}}$$

if  $\|v\|_E = \rho_o$ . (3.22) and (3.17) imply

$$F(v) > F(0)$$

provide that  $\|v\|_E = \rho_o$ . The proof is complete.  $\square$

It follows from (3.17) that

$$(3.23) \quad F(0) < \frac{1}{N} S^{\frac{N}{2}}.$$

We consider the problem

$$(3.24) \quad m := \min\{F(v) : v \in B_{\rho_o}\}.$$

**Lemma 3.3.** *Problem (3.1) has a nontrivial solution  $v_o \in B_{\rho_o}$ .*

*Proof.* By (3.23) we have

$$(3.25) \quad m < \frac{1}{N} S^{\frac{N}{2}}.$$

Let  $\{v_n\}$  be a minimizing sequence of (3.24). Since  $\|v_n\|_E \leq \rho_o$  we may assume

$$(3.26) \quad v_n \rightarrow v_o \text{ weakly in } E, \quad v_n \rightarrow v_o \text{ in } L^q(\Omega), \quad 2 \leq q < 2^*, \quad v_n \rightarrow v_o \text{ a.e. in } \Omega$$

as  $n \rightarrow \infty$ . The weak continuity of norm gives

$$(3.27) \quad \|v_o\|_E \leq \liminf_{n \rightarrow \infty} \|v_n\|_E \leq \rho_o.$$

By the Ekeland's variational principle, we may assume that

$$(3.28) \quad F(v_n) \rightarrow m, \quad F'(v_n) \rightarrow 0$$

as  $n \rightarrow \infty$ . Because of

$$(3.29) \quad F'(v_n) \rightarrow 0 \iff J'(v_n + t(v_n)\phi_1) \rightarrow 0$$

as  $n \rightarrow \infty$ , we have

$$(3.30) \quad \begin{aligned} & \frac{1}{2} \int_{\Omega} (|\nabla v_n|^2 - \lambda_1 v_n^2) dx - \frac{1}{2^*} \int_{\Omega} (v_n + t(v_n)\phi_1)_+^{2^*} dx - \int_{\Omega} f(v_n + t(v_n)\phi_1) dx \\ & = m + o(1) \end{aligned}$$

and

$$(3.31) \quad \int_{\Omega} (|\nabla v_n|^2 - \lambda_1 v_n^2) dx - \int_{\Omega} (v_n + t(v_n)\phi_1)_+^{2^*-1} v_n dx - \int_{\Omega} f v_n dx = o(1).$$

By the weak convergence we know that  $v_o$  satisfies

$$(3.32) \quad -\Delta v = \lambda_1 v + (v + t(v)\phi_1)_+^{2^*-1} + f,$$

and then

$$(3.33) \quad \int_{\Omega} (|\nabla v_o|^2 - \lambda_1 v_o^2 - (v_o + t(v_o)\phi_1)_+^{2^*-1} v_o - f v_o) dx = 0,$$

$$(3.34) \quad \int_{\Omega} [(v_o + t(v_o)\phi_1)_+^{2^*-1} \phi_1 + f \phi_1] dx = 0.$$

The proof will be complete if we may show  $v_o \not\equiv 0$ . First we claim that

$$(3.35) \quad \lim_{n \rightarrow \infty} t(v_n) = t(v_o).$$

If not, we would have  $\lim_n t(v_n) = t_1 \neq t(v_o)$ . By (3.6)

$$\int_{\Omega} (v_n + t(v_n)\phi_1)_+^{2^*-1} \phi_1 dx = - \int_{\Omega} f \phi_1 dx = \int_{\Omega} (v_o + t(v_o)\phi_1)_+^{2^*-1} \phi_1 dx,$$

it follows

$$\int_{\Omega} (v_o + t_1\phi_1)_+^{2^*-1} \phi_1 dx = \int_{\Omega} (v_o + t(v_o)\phi_1)_+^{2^*-1} \phi_1 dx$$

giving a contradiction. Letting  $w_n = v_n - v_o$ . By (3.30), (3.31) and Brézis-Lieb Lemma, we obtain

$$(3.36) \quad \begin{aligned} & \frac{1}{2} \int_{\Omega} |\nabla w_n|^2 dx - \frac{1}{2^*} \int_{\Omega} (w_n)_+^{2^*} dx + \frac{1}{2} \int_{\Omega} (|\nabla v_o|^2 - \lambda_1 v_o^2) dx \\ & - \frac{1}{2^*} \int_{\Omega} (v_o + t(v_o)\phi_1)_+^{2^*} dx - \int_{\Omega} f(v_o + t(v_o)\phi_1) dx = m + o(1), \end{aligned}$$

i.e.

$$(3.37) \quad F(v_o) + \frac{1}{2} \int_{\Omega} |\nabla w_n|^2 dx - \frac{1}{2^*} \int_{\Omega} (w_n)_+^{2^*} dx = m + o(1).$$

Similarly by (3.31), (3.34) and Brézis-Lieb Lemma we deduce

$$\begin{aligned} & \int_{\Omega} |\nabla w_n|^2 dx - \int_{\Omega} (w_n)_+^{2^*} dx - \int_{\Omega} (v_o + t(v_o)\phi_1)_+^{2^*} dx \\ & + \int_{\Omega} (|\nabla v_o|^2 - \lambda_1 v_o^2) dx - \int_{\Omega} f(v_o + t(v_o)\phi_1) dx = o(1), \end{aligned}$$

namely

$$(3.38) \quad \int_{\Omega} |\nabla w_n|^2 dx - \int_{\Omega} (w_n)_+^{2^*} dx = o(1).$$

Let  $\lim_{n \rightarrow \infty} \int_{\Omega} |\nabla w_n|^2 dx = k \geq 0$ . If  $k = 0$ , we have done. If  $k > 0$ , by the Sobolev inequality

$$(3.39) \quad \int_{\Omega} |\nabla w_n|^2 dx \geq S \left( \int_{\Omega} (w_n)_+^{2^*} dx \right)^{2/2^*}.$$

Taking the limit in (3.39) we obtain by (3.38) and (3.39) that

$$k \geq S k^{\frac{N-2}{N}}$$

i.e.

$$(3.40) \quad k \geq S^{\frac{N}{2}}.$$

It yields by (3.37), (3.39) and (3.40)

$$\frac{1}{N}S^{\frac{N}{2}} > m \geq F(v_o) + \frac{1}{N}k \geq F(v_o) + \frac{1}{N}S^{\frac{N}{2}}.$$

So  $F(v_o) < 0$  and  $v_o \neq 0$ . Since  $F(v) \geq \alpha > 0$  if  $\|v\|_E = \rho_o$ , we have  $v_o \in B_{\rho_o}$ .

The proof is complete.

□

#### 4. BIFURCATIONS AT $\lambda = \lambda_k$

In this section we discuss the bifurcation of the set of solutions of (1.1). Let  $u_t(\lambda) = u_t$  be the negative solution obtained in Section 2. If  $f = t\phi_1 + h$  and  $h \in \ker(-\Delta - \lambda)^\perp$ ,  $u_t(\lambda)$  is well defined for all  $\lambda \neq \lambda_1$ . In the case  $\lambda = \lambda_k, k \neq 1$ , the set of solutions of (1.1) bifurcating from  $(\lambda_k, u_t(\lambda_k))$  is equivalent to the set of solutions of (2.4) bifurcating from  $(\lambda_k, 0)$ . Let

$$E^- = \text{span}\{\phi_1, \dots, \phi_k\}, \quad E^+ = (E^-)^\perp.$$

Now we state a bifurcation result.

**Proposition 4.1.** *Every eigenvalue  $\lambda_k$  of  $-\Delta$  gives rise to a bifurcation point of  $(\lambda_k, 0)$  of (2.4). As a result, we obtain Theorem 1.3.*

*Proof.* The conclusion follows from an abstract bifurcation theorem due to Böhme[5] and Marino[21], see also Theorem 11.4 in [24]. Let  $\chi(\xi) \in C^\infty(\mathbb{R}, \mathbb{R})$  satisfy  $\chi(\xi) = 1$  for  $|\xi| \leq 1$ ,  $\chi(\xi) = 0$  and  $|\xi| \geq 2$ , and  $0 \leq \chi(\xi) \leq 1$  for all  $\xi$ . Define

$$g(\lambda, \xi) = \chi(\xi)(\xi + u_t(\lambda))_+^{2^*-1} + (1 - \chi(\xi)).$$

Then  $g \in C^1$ , and  $g(\lambda, \xi) = o(|\xi|)$  for  $\lambda$  bounded. Set

$$\Phi(v) = \frac{1}{2} \int_{\Omega} |\nabla v|^2 dx - \int_{\Omega} G(\lambda, v) dx$$

with  $u \in E := W_o^{1,2}(\Omega)$ , where  $G(\lambda, v) = \int_0^v g(\lambda, t) dt$ . It is standard to show that  $\Phi \in C^2$ . A critical point  $u$  of  $\Phi$  on the manifold  $M := \{u \in E : \int_{\Omega} |u|^2 dx = r^2\}$  is a weak solution of

for some Lagrange multiplier  $\gamma$ . Define the operator  $L$  by

$$(Lv, \phi) = \int_{\Omega} \nabla v \nabla \phi \, dx$$

and  $H$  by

$$H(v)\phi = \int_{\Omega} g(\lambda, v)\phi \, dx$$

for  $\phi \in E$ . For any  $\nu$  satisfies  $2 < \nu < 2^*$  and  $\omega := \{x \in \Omega : v(x) \geq 2\}$  with  $v \in E$ ,

we have

$$\int_{\Omega} |v|^{\nu} \, dx \geq 2^{\nu} \text{meas}\omega.$$

Hence

$$|H(v)\phi| \leq \int_{\Omega/\omega} |v|^{2^*-1} |\phi| \, dx + \int_{\omega} |\phi| \, dx \leq C \|v\|^{\nu} \|\phi\|_E.$$

It concludes

$$\|H(v)\| = o(\|v\|).$$

So by Theorem 11.4 in [24], each eigenvalue of  $-\Delta$  provides a bifurcation point of

$$(4.1) \quad -\Delta v - g(\lambda, v) = \lambda v.$$

Since  $g(\lambda, v) = o(|v|)$  and  $\lambda$  is bounded, it follows from (4.1) that

$$\|v\|_E \leq C \|v\|_{L^2(\Omega)} = Cr.$$

Arguments from elliptic regularity theory [6] show if  $r$  is small enough,  $\|v\|_{L^\infty(\Omega)} < 1$  and  $g(\lambda, v) = (v + u_t(\lambda))_+^{2^*-1}$ . The proof is complete.

□

Next, we show that the bifurcation branch bends locally to the left.

**Proposition 4.2.** *If  $(\lambda, v(\lambda)), v(\lambda) \neq 0$ , is a solution of (2.4) such that  $\lambda \rightarrow \lambda_k, k \neq 1, v(\lambda) \rightarrow 0$ , then  $\lambda < \lambda_k$ . Consequently, if  $h \in \ker(-\Delta - \lambda_k)^\perp$  and  $(\lambda, u(\lambda)), u(\lambda) \neq 0$ , is a solution of (1.1) such  $\lambda \rightarrow \lambda_k, k \neq 1$  and  $u(\lambda) \rightarrow u_t(\lambda_k)$ , then  $\lambda < \lambda_k$ .*

*Proof.* Let  $u = v + w$  be a solution of (2.4) with  $v \in E^-$  and  $w \in E^+$ . Multiplying (2.4) by  $w - v$  and integrating by part, we obtain

$$(4.2) \quad \begin{aligned} & \int_{\mathbb{R}^N} (|\nabla w|^2 - |\nabla v|^2) \, dx \\ &= \int_{\mathbb{R}^N} [\lambda u + (u + u_t(\lambda))_+^{2^*-1}] \, dx \\ &= \int_{\mathbb{R}^N} [\lambda(w^2 - v^2) + (v + w + u_t(\lambda))_+^{2^*-1}] \, dx. \end{aligned}$$

It follows

$$(4.3) \quad \left(1 - \frac{\lambda}{\lambda_{k+1}}\right) \int_{\mathbb{R}^N} |\nabla w|^2 dx - \left(1 - \frac{\lambda}{\lambda_k}\right) \int_{\mathbb{R}^N} |\nabla v|^2 dx \leq \int_{\mathbb{R}^N} (v+w+u_t(\lambda))_+^{2^*-1} (w-v) dx.$$

By the convexity of the function  $(v+w+u_t(\lambda))_+^{2^*-1}$  and since  $u_t$  is negative

$$(4.4) \quad \begin{aligned} & \int_{\mathbb{R}^N} (v+w+u_t(\lambda))_+^{2^*-1} (w-v) dx = \int_{\mathbb{R}^N} (v+w+u_t(\lambda))_+^{2^*-1} (2w-u) dx \\ & \leq \int_{\mathbb{R}^N} (2w+u_t(\lambda))_+^{2^*} dx - \int_{\mathbb{R}^N} (u+u_t(\lambda))_+^{2^*} dx \\ & \leq \int_{\mathbb{R}^N} (2w+u_t(\lambda))_+^{2^*} dx \leq \int_{\mathbb{R}^N} |2w|^{2^*} dx \leq C \|w\|_E^{2^*}. \end{aligned}$$

(4.3) and (4.4) imply

$$(4.5) \quad \left[\left(1 - \frac{\lambda}{\lambda_{k+1}}\right) - C \|w\|_E^{2^*-2}\right] \|w\|_E^2 - \left(1 - \frac{\lambda}{\lambda_k}\right) \|v\|_E^2 \leq 0.$$

Suppose by contradiction that  $\lambda \geq \lambda_k$ . Since  $\frac{\lambda}{\lambda_k} - 1 > 0$  and  $u = v + w \neq 0$  we must have  $w \neq 0$ . Hence

$$(4.6) \quad \left(1 - \frac{\lambda}{\lambda_{k+1}}\right) \leq C \|w\|_E^{2^*-2} \leq C \|u\|_E^{2^*-2}.$$

It yields a contradiction when we let  $\lambda \rightarrow \lambda_k$ .  $\square$

## REFERENCES

- [1] H.Amann and P.Hess, *A multiplicity result for a class of elliptic boundary value problems*, Proc. Royal Soc. Edinburgh **84A** (1979), 145-151.
- [2] A.Ambrosetti and G.Prodi, *On the inversion of some differentiable mappings with singularities between Banach spaces*, Ann. Math. Pura Appl. **93** (1972), 231-247.
- [3] H.Berestycki, *Le nombre de solutions de certains problèmes semi-linéaire elliptiques*, J.Funct. Anal. **40** (1981), 1-29.
- [4] M.S.Berger and E.Podolak, *On the solutions of a nonlinear Dirichlet problem*, Indiana Univ. Math. J. **24** (1975), 837-846.
- [5] R.Böhme, *Die Lösung der Verzweigungsgleichungen für nichtlineare Eigenwertprobleme*, Math.Z. **127** (1972), 105-126.
- [6] H.Brézis and T.Kato, *Remarks on the Schrödinger operator with singular complex potentials*, J.Math.Pures et Appl. **58** (1978), 137-151.
- [7] H.Brézis and E.Lieb, *A relation between pointwise convergence of functions and convergence of integrals*, Proc. Amer. Math. Soc. **88** (1983), 486-490.
- [8] H.Brézis and L.Nirenberg, *Positive solutions of nonlinear elliptic equations involving critical Sobolev exponents*, Comm. Pure Appl. Math. **24** (1983), 437-477.
- [9] A.Castro, *Metodos Variacionales y Analisis Funcional no Lineal*, X Coloquio Colombiano de Matematicos 1980.
- [10] A.Castro, *Reduction via minimax*, Lecture Notes in Math. **957**, Springer.
- [11] J.Chabrowski and Yang Jianfu, *Existence theorems for the Schrödinger equation involving a critical Sobolev exponent*, ZAMP **48** (1998), 276-293.
- [12] J.N.Dancer, *On the ranges of certain weakly nonlinear elliptic partial differential equations*, J. Math. Pures et Appl. **57** (1978), 251-266.



- [13] Peng Yinbin, *On the superlinear Ambrosetti-Prodi problem involving critical Sobolev exponents*, Nonlinear Anal. TMA **17(12)** (1991), 1111-1124.
- [14] P.G.de Figueiredo, *The Ekeland variational principle with applications and detours*, Tate IFR Lectures on Math.and Phys. **81**, Springer 1989.
- [15] P.G.de Figueiredo, *On superlinear elliptic problems with nonlinearities interacting only with higher eigenvalues*, Rocky Mountain J. Math. **18** (1988), 287 - 303.
- [16] P.G.de Figueiredo, *Lectures on boundary value problems of Ambrosetti-Prodi type*, Atas do 12<sup>o</sup> Seminario Brasileiro de Análise. São Paulo, 1980.
- [17] J.Fučík, *Solvability of nonlinear equations and boundary value problems*, D. Reidel Publ.Co., Dordrecht 1980.
- [18] L.Kazdan and F.W.Warner, *Remarks on some quasilinear elliptic equations*, Comm. Pure and Appl. Math. **XXVIII** (1975), 567-597.
- [19] K.Küpper and C.A.Stuart, *Bifurcation into gaps in the essential spectrum*, J.Reine Angew. Math. **409** (1990), 1-34.
- [20] C.Lazer and P.J.McKenna, *Nonlinear perturbations of linear elliptic boundary value problems at resonance*, J.Math.Mech. **19** (1973), 63-72.
- [21] A.Marino, *La biforcazione nel caso variazionale*, 132 Confer.Sem.Mat.Univ.Bari 1977.
- [22] Mawhin and M.Willem, *Critical Point Theory and Hamiltonian Systems*, Springer-Verlag, 1989.
- [23] M.Mitrović and D.Zubrinić, *Fundamentals of Applied Functional Analysis*, Pitman Monographs and Surveys in Pure and Applied Mathematics No: 91. Addison Wesley Longman, 1998.
- [24] H.Rabinowitz, *Minimax methods in critical point theory with applications to differential equations*, 65 AMS Conf. Ser. Math. 1986.
- [25] R.Ruf and P.N.Srikanth, *Multiplicity results for superlinear elliptic problems with partial interference with spectrum*, J.Math.Anal.Appl. **118** (1986), 15-23.
- [26] T.Tarantello, *On nonhomogeneous elliptic equations involving critical Sobolev exponent*, Ann. I.H.P.Anal. Non linéaire **9** (1992), 281-304.
- [27] M.Willem, *Minimax Theorems*, Progress in Nonlinear Differential Equations and their Applications, 24 Birkhäuser, Boston, Inc Boston MA. 1996.