

# Some Nonlinear Perturbation of the Quasilinear hyperbolic Equation\*

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## Abstract

In this paper we prove existence and uniqueness of global regular solutions of a mixed problem for the perturbed of extensional vibrations of a thin rod equation or torsional vibrations of thin rod.

## 1 Introduction

Let us consider a Lipschitzian perturbation of vibrations of a thin rod operator (see. Love [1])

$$Ku = \frac{\partial^2 u}{\partial t^2} - \Delta u - M \left( \int_{\Omega} |\nabla u(x, t)|^2 dx \right) \Delta u_{tt},$$

given by  $Ku + F(u)$ , where  $F$  is a Lipschitzian function which satisfies some appropriate conditions fixed in Section 2 and  $M \in C^1([0, +\infty))$ . We obtain the nonlinear mixed problem:

$$\begin{aligned} \frac{\partial^2 u}{\partial t^2} - \Delta u - M \left( \int_{\Omega} |\nabla u(x, t)|^2 dx \right) \Delta u_{tt} + F(u) &= f \text{ in } Q, \\ u &= 0 \text{ on } \Sigma, \\ u(x, 0) &= u_0(x), \quad u_t(x, 0) = u_1(x) \text{ in } \Omega. \end{aligned} \tag{1}$$

In (1)  $\Omega$  is a bounded open set of  $R^n$  with smooth boundary  $\Gamma$ ,  $Q$  is the cylinder  $\Omega \times (0, T)$ ,  $T$  is a positive real number. The lateral boundary of  $Q$

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is represented by  $\Sigma$ , i.e.  $\Sigma = \Gamma \times (0, T)$ . By  $M(\lambda)$  we denote a real function defined on the positive real numbers with  $M(\lambda) \geq \rho > 0$  for all  $\lambda \geq 0$ , for some  $\rho > 0$ ;  $\nabla u$  is the gradient of  $u$  and  $\Delta$  is the Laplace operator.

Solving equation (1) is important by the following reasons: equation (1) with  $M(\lambda) = 1$  and  $F(u) = 0$  arises in the study of the extensional vibrations of thin rods; many authors studied this equation when  $M(\lambda) = 1$ ; among them we mention Pereira [12] and Ferreira and Pereira [13]. Besides, equation (1) with  $F(u) = 0$ ,  $M(\lambda) = \lambda_0$ ,  $\lambda_0 = \int_{\Omega} \phi^2 dx$ , where  $\phi$  is the torsion-function, also appears in the study of the torsional vibrations of thin rods (see. Love [1]). Finally, the function  $M(\lambda)$  in (1) has its motivation in the mathematical description of the vibrations of an elastic stretched string

$$\frac{\partial^2 u}{\partial t^2} - M \left( \int_{\Omega} |\nabla u(x, t)|^2 dx \right) \Delta u = 0, \quad (2)$$

which for  $M(\lambda) \geq \rho > 0$ , was studied by Pohozaev [16], Nishihara [15] and Lions [11]. When  $M(\lambda) \geq 0$  was treated by Arosio and Spagnolo [8], Ebihara, Medeiros and Miranda [4], Yamada [6], Matos [14]. When  $M(\lambda)$  is a constant  $C^2$ , (2) becomes a perturbation of the d'Alembert operator

$$\square u = \frac{\partial^2 u}{\partial t^2} - C^2 \Delta u$$

(see Lions [11]). In [2], Strauss studied a nonlinear perturbation of this operator of the type:

$$\left\{ \begin{array}{l} \square u + F(u) = f \text{ in } Q \\ u = 0 \text{ on } \Sigma \\ u(0) = u_0, u'(0) = u_1 \text{ in } \Omega \end{array} \right. \quad (3)$$

where  $F : R \rightarrow R$  is continuous and  $sF(s) \geq 0$  for all  $s \in R$ .

For others perturbations of Kirchhoff-Carrier operator, see Jörgens [9], Hosoya and Yamada [5], D'Ancona and Spagnolo [7].

In Ebihara and Pereira [3] it was proved that there exists only one classical solution for a quasilinear model, given by following initial-boundary value problem:

$$\left\{ \begin{array}{l} \frac{\partial^2 u}{\partial t^2} - \Delta u - M \left( \int_{\Omega} |\nabla u(x, t)|^2 dx \right) \Delta u_{tt} = f \text{ in } Q \\ u = 0 \text{ on } \Sigma \\ u(0) = u_0, u'(0) = u_1 \text{ in } \Omega \end{array} \right. \quad (4)$$

when the following hypotheses hold:

**(H.0)**  $M(\lambda) \in C^1([0, +\infty))$ , and there exist positive constants  $\alpha, \rho$  such that

$$M(\lambda) \geq \alpha\sqrt{\lambda} + \rho, \text{ for all } \lambda \in [0, +\infty).$$

**(H.1)** There exists a nonnegative function  $\beta(\lambda)$  satisfying:

$$\left| \frac{d}{d\lambda} M(\lambda) \right| \sqrt{\lambda} \leq \beta(\lambda) M(\lambda), \text{ for all } \lambda \in [0, +\infty).$$

**(H.2)** The initial conditions are such that:

$$u_0, u_1 \in D(A^l), \quad l \geq 2$$

and

$$f, \frac{df}{dt} \in C(0, T; D(A^{\frac{l-1}{2}})), \quad l \geq 2$$

where  $A = -\Delta$  and by  $D(A^s)$  we are denoting the domain of the operator  $A^s$ .

Motivated by Hosoya and Yamada [5], Strauss [2] and Ebihara and Pereira [3], we shall investigate in this paper the following perturbation of the vibrations model,

$$\begin{cases} Ku + F(u) = f \text{ in } Q \\ u = 0 \text{ on } \Sigma \\ u(0) = u_0, \quad u_1(0) = u_1 \text{ in } \Omega \end{cases} \quad (5)$$

with  $F$  Lipschitzian.

Our main objective is to give a complete, clear and short proof for the existence and uniqueness of global regular solutions to the problem in question. For that, we use the Faedo-Galerkin's method associated with a compactness argument and some technical ideas. In our proof of the existence and uniqueness of global regular solutions to (5) we assume **(H.0)** and substitute **(H.2)** by  $f \in C([0, T]; H_0^1(\Omega))$  and  $u_0, u_1 \in D(A^{1/2})$  and **(H.1)** becomes unnecessary.

## 2 Notations and Main Result

For the Hilbert space  $L^2(\Omega)$  we denote its inner product and norm, respectively, by  $(\cdot, \cdot)$  and  $|\cdot|$ , defined by:

$$(u, v) = \int_{\Omega} uv dx ; \quad |u|^2 = \int_{\Omega} |u|^2 dx.$$

By  $H^m(\Omega)$  we represent the Sobolev space on the bounded open set  $\Omega$  of  $R^n$ .  $H_0^m(\Omega)$  is the closure of  $C_0^\infty(\Omega)$  in  $H^m(\Omega)$ . In particular,  $H^1(\Omega)$  has inner product  $((\cdot, \cdot))$  and norm  $\|\cdot\|$  given by

$$((u, v)) = \int_{\Omega} uv dx + \int_{\Omega} \nabla u \cdot \nabla v dx$$

and

$$\|v\|^2 = \int_{\Omega} v^2 dx + \int_{\Omega} |\nabla v|^2 dx.$$

In  $H_0^1(\Omega)$  we consider the equivalent norm

$$\|v\|^2 = \int_{\Omega} |\nabla v|^2 dx$$

and the inner product

$$((u, v)) = \int_{\Omega} \nabla u \cdot \nabla v dx$$

for all  $u, v \in H_0^1(\Omega)$ . We also observe that if  $\Gamma$  is of class  $C^{1,1}$ , then the norms of  $\|v\|_{H^2}$  and  $|\Delta v|$  are equivalent for  $v \in H_0^1(\Omega) \cap H_0^2(\Omega)$ .

Suppose that the functions  $M(\lambda)$  and  $F(u)$  satisfy:

**(A.1)**  $M \in C^1([0, +\infty))$  and there exist positive constants  $\alpha, \rho$  such that the following inequality is valid:

$$M(\lambda) \geq \alpha\sqrt{\lambda} + \rho, \text{ for all } \lambda \in [0, +\infty).$$

**(A.2)**  $F : R \rightarrow R$  is a Lipschitzian function such that

$$sF(s) \geq 0 \text{ for all } s \in R.$$

We have the following result:

**THEOREM 1.** Suppose  $u_0, u_1 \in H_0^1(\Omega) \cap H^2(\Omega)$ ,  $f \in C([0, T]; H_0^1(\Omega))$  and assumptions **(A.1)**-**(A.2)** hold. Then there exists a unique function  $u : [0, T] \rightarrow L^2(\Omega)$ , satisfying:

$$u \in C^1([0, T]; D(A^\alpha)) \text{ for all } 0 \leq \alpha < 1 \tag{6}$$

$$u' \in L^\infty(0, T; H_0^1(\Omega) \cap H^2(\Omega)) \tag{7}$$

$$u'' \in L^\infty(0, T; H_0^1(\Omega) \cap H^2(\Omega)) \tag{8}$$

$$Ku + F(u) = f \text{ in } L^\infty(0, T; L^2(\Omega)) \tag{9}$$

$$u(0) = u_0 \quad (10)$$

$$u'(0) = u_1 \quad (11)$$

$$u = 0 \text{ on } \Sigma. \quad (12)$$

As it is well known, the Laplace operator with homogeneous boundary conditions has a set of eigenvectors such that the subspace spanned by them is dense in  $L^2(\Omega)$ , in  $H_0^1(\Omega)$  and in  $H_0^1(\Omega) \cap H_0^2(\Omega)$ . We may consider these eigenvectors to be orthonormal in  $L^2(\Omega)$  and orthogonal in  $H_0^1(\Omega)$  or  $H_0^1(\Omega) \cap H_0^2(\Omega)$ . Moreover, by the regularity of the boundary  $\Gamma$ , we have  $w_j(x) \in H^2(\Omega)$ .

Let us denote by  $V_m = [w_1, \dots, w_m]$  the subspace of  $H_0^1(\Omega)$  generated by the first  $m$  eigenvectors  $w_j$  ( $m \geq 1$ ).

Let

$$u_m(t) = \sum_{j=1}^m g_{jm}(t)w_j(x) \in V_m$$

be a solution of the system:

$$(Ku_m(t), v) + (F(u_m(t), v) = (f(t), v), \quad \forall v \in V_m, \quad (13)$$

$$u_m(0) = u_{0m} \rightarrow u_0 \text{ strongly in } H_0^1(\Omega) \cap H^2(\Omega) \text{ as } m \rightarrow \infty, \quad (14)$$

$$u'_m(0) = u_{1m} \rightarrow u_1 \text{ strongly in } H_0^1(\Omega) \cap H^2(\Omega) \text{ as } m \rightarrow \infty, \quad (15)$$

where

$$u_{0m} = \sum_{j=1}^m (u_0, w_j)w_j, \quad u_{1m} = \sum_{j=1}^m (u_1, w_j)w_j.$$

Setting  $\nu = w_j$  for  $j = 1, \dots, m$  we obtain a system of ordinary differential equations to which the classical Peano's theorem is applicable. Therefore this system has a local solution  $u_m(t)$  on  $[0, T_m)$ ,  $0 < T_m \leq T$ . The extension to  $[0, T]$  is a consequence of the estimate (i) given in the next Section.

To prove our results, we will use the following result due to Simon [17]

**Lemma 2.1** *Let  $X \subset E \subset Y$  Banach spaces, the imbedding  $X \rightarrow E$  being compact. Then the following imbeddings are compact:*

$$(i) L^q(0, T; X) \cap \{\varphi : \frac{\partial \varphi}{\partial t} \in L^1(0, T; Y)\} \rightarrow L^q(0, T; E) \text{ if } 1 \leq q \leq \infty,$$

$$(ii) L^\infty(0, T; X) \cap \{\varphi : \frac{\partial \varphi}{\partial t} \in L^r(0, T; Y)\} \rightarrow C([0, t]; E) \text{ if } 1 < r \leq \infty.$$

and (see Temam [18])

**Lemma 2.2** *Let  $X \subset E \subset X^*$  Banach spaces. If  $u \in L^2(0, T; X)$  and  $u_t \in L^2(0, T; X^*)$  then  $u \in C([0, T]; E)$  a.e.*

### 3 A Priori Estimates

**Estimate (i).** Setting  $v = u_m''(t)$  in (13), we obtain

$$\begin{aligned} |u_m''(t)|^2 + a(u_m(t), u_m''(t)) + M(a(u_m(t)))a(u_m''(t), u_m''(t)) \\ + (F(u_m(t)), u_m''(t)) = (f(t), u_m''(t)). \end{aligned} \quad (16)$$

where

$$a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v dx \text{ and } a(u) = \int_{\Omega} |\nabla u|^2 dx.$$

Let us define  $\mu(t) = M(a(u_m(t))) \geq \rho > 0$ . Dividing both sides of (16) by  $\mu(t)$ , applying Cauchy-Schwartz's inequality and using the hypothesis **(A.1)**, we get

$$\begin{aligned} \frac{1}{\mu(t)} |u_m''(t)|^2 + |\nabla u_m''(t)|^2 \leq & \frac{\|u_m(t)\| |\nabla u_m''(t)|}{\alpha \|u_m(t)\|} + \frac{|f(t)| |u_m''(t)|}{\rho} \\ & + \frac{|F(u_m(t))| |u_m''(t)|}{\alpha \|u_m(t)\|}. \end{aligned} \quad (17)$$

Since  $F$  is Lipschitzian and satisfies  $sF(s) \geq 0$  for all  $s \in R$  (thus,  $F(0) = 0$ ) we obtain

$$|F(u_m(t))| \leq C \left( \int_{\Omega} |u_m(t)|^2 dx \right)^{\frac{1}{2}} = C |u_m(t)|. \quad (18)$$

Using the Poincaré's inequality in (18), we have

$$|F(u_m(t))| \leq C |\nabla u_m(t)|, \quad C > 0. \quad (19)$$

Applying (19) in (17), yields

$$\begin{aligned} \frac{1}{\mu(t)} |u_m''(t)|^2 + |\nabla u_m''(t)|^2 \leq & \frac{1}{\alpha^2} + \frac{1}{4} |\nabla u_m''(t)|^2 + \frac{C^2}{\rho^2} |\nabla f(t)|^2 \\ & + \frac{1}{4} |\nabla u_m''(t)|^2 + \frac{C^2}{\alpha^2} + \frac{1}{4} |\nabla u_m''(t)|^2 \end{aligned} \quad (20)$$

where  $C$  is a positive constant such that  $|f(t)| \leq C |\nabla f(t)|$  and  $|u_m''(t)| \leq C |\nabla u_m''(t)|$ .

Then, from the above estimate and  $f \in C([0, T]; H_0^1(\Omega))$ , we get

$$|\nabla u_m''(t)|^2 \leq 4 \left( \frac{1}{\alpha^2} + \frac{C^2}{\alpha^2} + \frac{C^2}{\rho^2} \max_{0 \leq t \leq T} |\nabla f(t)|^2 \right) < C_1 \quad (21)$$

where  $C_1 > 0$  is independent of  $t$  and  $m$ .

Thus,

$$(u_m'') \text{ is bounded in } L^\infty(0, T_m; H_0^1(\Omega)). \quad (22)$$

By the Fundamental Theorem of Calculus, we have

$$(u_m') \text{ is bounded in } L^\infty(0, T_m; H_0^1(\Omega)) \quad (23)$$

and

$$(u_m) \text{ is bounded in } L^\infty(0, T_m, H_0^1(\Omega)). \quad (24)$$

The above estimates permit us to extend  $u_m(t)$  to all interval  $[0, T]$ .

**Estimate (ii).** Since we are working with spectral basis, we have  $-\Delta u_m'(t) \in V_m$ . Then we can take  $v = -2\Delta u_m'(t)$  in (13) to obtain

$$\begin{aligned} \frac{d}{dt} [|\nabla u_m'(t)|^2 + |\Delta u_m(t)|^2] + M(a(u_m(t))) \frac{d}{dt} |\Delta u_m'(t)|^2 \\ + 2(\nabla F(u_m(t)), \nabla u_m'(t)) = 2(\nabla f(t), \nabla u_m'(t)). \end{aligned}$$

Setting  $\mu(t) = M(a(u_m(t))) \geq \rho > 0$ , dividing both sides of the last equality by  $\mu(t)$ , we obtain

$$\begin{aligned} \frac{1}{\mu(t)} \frac{d}{dt} [|\nabla u_m'(t)|^2 + |\Delta u_m(t)|^2] + \frac{d}{dt} |\Delta u_m'(t)|^2 \\ = \frac{-2}{\mu(t)} (\nabla F(u_m(t)), \nabla u_m'(t)) + \frac{2}{\mu(t)} (\nabla f(t), \nabla u_m'(t)). \end{aligned} \quad (25)$$

Defining

$$h_m(t) = \frac{1}{\mu(t)} [|\nabla u_m'(t)|^2 + |\Delta u_m(t)|^2] + |\Delta u_m'(t)|^2 \quad (26)$$

we have

$$\begin{aligned} h_m'(t) &= \frac{1}{\mu(t)} \frac{d}{dt} [|\nabla u_m'(t)|^2 + |\Delta u_m(t)|^2] + \frac{d}{dt} |\Delta u_m'(t)|^2 \\ &\quad - \frac{\mu'(t)}{\mu(t)^2} [|\nabla u_m'(t)|^2 + |\Delta u_m(t)|^2]. \end{aligned} \quad (27)$$

Using (25) in (27), we obtain

$$\begin{aligned} h'_m(t) &\leq \frac{2}{\mu(t)} |\nabla F(u_m(t))| |\nabla u'_m(t)| + \frac{2}{\mu(t)} |\nabla f(t)| |\nabla u'_m(t)| \\ &\quad + \frac{|\mu'(t)|}{\mu(t)^2} [|\nabla u'_m(t)|^2 + |\Delta u_m(t)|^2]. \end{aligned}$$

We observe that

$$|\nabla F(u_m(t))| = |F'(u_m(t)) \nabla u_m(t)| \leq C |\nabla u_m(t)|$$

because  $F$  is Lipschitzian.

Consequently,

$$\begin{aligned} h'_m(t) &\leq \frac{2C}{\mu(t)} |\nabla u_m(t)| |\nabla u'_m(t)| + \frac{2}{\mu(t)} |\nabla f(t)| |\nabla u'_m(t)| \quad (28) \\ &\quad + \frac{|\mu'(t)|}{\mu(t)^2} [|\nabla u'_m(t)|^2 + |\Delta u_m(t)|^2]. \end{aligned}$$

We also observe that

$$\mu'(t) = M'(a(u_m(t))) 2a(u_m(t), u'_m(t))$$

and

$$\begin{aligned} |M'(a(u_m(t)))| &\leq C_2 \text{ if } 0 \leq t \leq T, \ 0 \leq a(u_m(t)) \leq C_1 \\ |a(u_m(t), u'_m(t))| &\leq |\nabla u_m(t)| |\nabla u'_m(t)|. \end{aligned}$$

Then

$$|\mu'(t)| \leq C_3,$$

where  $C_3$  is a positive constant.

From (28), we get

$$\begin{aligned} h'_m(t) &\leq \frac{2C}{\mu(t)} |\nabla u_m(t)| |\nabla u'_m(t)| + \frac{2}{\mu(t)} |\nabla f(t)| |\nabla u'_m(t)| \quad (29) \\ &\quad + \frac{C_3}{\mu(t)^2} [|\nabla u'_m(t)|^2 + |\Delta u_m(t)|^2]. \end{aligned}$$



Using the estimate (i), we get

$$h'(t) \leq \frac{C_4^2}{\mu(t)} + \frac{1}{\mu(t)} |\nabla u'_m(t)|^2 + \frac{1}{\mu(t)} |\nabla f(t)|^2 + \frac{1}{\mu(t)} |\nabla u'_m(t)|^2 \quad (30)$$

$$+ \frac{C_3}{\mu(t)} \frac{1}{\mu(t)} [|\nabla u'_m(t)|^2 + |\Delta u_m(t)|^2].$$

By the definition of  $h_m(t)$ ,

$$\frac{1}{\mu(t)} [|\nabla u'_m(t)|^2 + |\Delta u_m(t)|^2] \leq h_m(t),$$

and

$$\frac{1}{\mu(t)^{\frac{1}{2}}} |\Delta u_m(t)| \leq h_m(t)^{\frac{1}{2}}.$$

From (23), we have

$$\frac{1}{\mu(t)} |\nabla u'_m(t)|^2 \leq C_1,$$

where  $C_1$  is positive constant.

Hence

$$h'_m(t) \leq \frac{C_4^2}{\rho} + 2h_m(t) + \frac{1}{\rho} \max_{0 \leq t \leq T} |\nabla f(t)|^2 + \frac{C_3}{\rho^{\frac{1}{2}}} h_m(t).$$

Thus,

$$\frac{d}{dt} h_m(t) \leq C_5 (1 + h_m(t)), \quad 0 \leq t \leq T, \quad (31)$$

where

$$C_5 = \max \left\{ \frac{C_4^2}{\rho} + \frac{1}{\rho} \max_{0 \leq t \leq T} |\nabla f(t)|^2, 2, \frac{C_3}{\rho^{\frac{1}{2}}} \right\} > 0.$$

From (31), we can easily see that, there exists a positive constant  $C_6$  such that for all  $t \in [0, T]$

$$h_m(t) \leq C_6 \text{ in } [0, T].$$

This implies that

$$\frac{1}{\mu(t)} [|\nabla u'_m(t)|^2 + |\Delta u_m(t)|^2] + |\Delta u'_m(t)|^2 < C_6.$$

Therefore,

$$\left[ |\nabla u'_m(t)|^2 + |\Delta u_m(t)|^2 \right] + \rho |\Delta u'_m(t)|^2 < C_7, \quad 0 \leq t \leq T \quad (32)$$

where  $C_7 > 0$  is a constant independent of  $t$  and  $m$ .

Whence,

$$(u_m) \text{ is bounded in } L^\infty(0, T; H^2(\Omega)), \quad (33)$$

$$(u'_m) \text{ is bounded in } L^\infty(0, T; H^2(\Omega)). \quad (34)$$

**Estimate (iii).** Setting  $v = -\Delta u''_m(t)$  in (13), we obtain

$$\begin{aligned} |\nabla u''_m(t)|^2 + (\Delta u_m(t), \Delta u''_m(t)) + M(a(u_m(t))) |\Delta u''_m(t)|^2 \\ + (F(u_m(t)), -\Delta u''_m(t)) = (f(t), -\Delta u''_m(t)). \end{aligned}$$

Let  $\mu(t) = M(a(u_m(t))) \geq \rho > 0$ . Dividing both sides of the last equality by  $\mu(t)$ , applying Cauchy-Schwartz's inequality and using hypothesis **(A.1)**, we get

$$\begin{aligned} \frac{1}{\mu(t)} |\nabla u''_m(t)|^2 + |\Delta u''_m(t)|^2 \leq \frac{1}{\rho} |\Delta u''_m(t)| + \frac{1}{\rho} |f(t)| |\Delta u''_m(t)| \quad (35) \\ + \frac{|F(u_m(t))| |\Delta u''_m(t)|}{\alpha \|u_m(t)\|}. \end{aligned}$$

Using (19),  $|f(t)| \leq C |\nabla f(t)|$  and recalling that  $|\Delta u'_m(t)|$  is bounded by the second estimate in (35), we obtain that

$$\begin{aligned} \frac{1}{\mu(t)} |\nabla u''_m(t)|^2 + |\Delta u''_m(t)|^2 \leq \frac{C_7^2}{\rho^2} + \frac{1}{4} |\Delta u''_m(t)|^2 + \frac{C^2}{\rho^2} |\nabla f(t)|^2 \quad (36) \\ + \frac{1}{4} |\Delta u''_m(t)|^2 + \frac{C^2}{\alpha^2} + \frac{1}{4} |\Delta u''_m(t)|^2. \end{aligned}$$

In particular we have

$$|\Delta u''_m(t)|^2 \leq 4 \left( \frac{C_7^2}{\rho^2} + \frac{C^2}{\alpha^2} + \frac{C^2}{\rho^2} \max_{0 \leq t \leq T} |\nabla f(t)|^2 \right) < C_8$$

where  $C_8 > 0$  is a constant independent of  $t$  and  $m$ .

Hence,

$$(u''_m) \text{ is bounded in } L^\infty(0, T; H^2(\Omega)). \quad (37)$$

## 4 Proof of Theorem

By estimates (22), (23), (24), (33), (34) and (37) there exists a subsequence of  $(u_m)_{m \in \mathbb{N}}$  which we still denote by  $(u_m)_{m \in \mathbb{N}}$ , and a function  $u$  such that:

$$u_m \rightarrow u \text{ weak - star in } L^\infty(0, T; H_0^1(\Omega) \cap H^2(\Omega)), \quad (38)$$

$$u'_m \rightarrow u' \text{ weak - star in } L^\infty(0, T; H_0^1(\Omega) \cap H^2(\Omega)), \quad (39)$$

$$u''_m \rightarrow u'' \text{ weak - star in } L^\infty(0, T; H_0^1(\Omega) \cap H^2(\Omega)). \quad (40)$$

From (38), (39) and (40), via Lemma 1, there exists a subsequence (still represented by  $(u_m)_{m \in \mathbb{N}}$ ) such that:

$$u_m \rightarrow u \text{ strongly in } C(0, T; D(A^\alpha)). \quad (41)$$

and

$$u'_m \rightarrow u' \text{ strongly in } C(0, T; D(A^\alpha)). \quad (42)$$

for all  $\alpha \in [\frac{1}{2}, 1)$

As a consequence,

$$M(a(u_m(t))) \rightarrow M(a(u(t))) \text{ strongly in } C(0, T; D(A^\alpha)). \quad (43)$$

Now, by (40), we have

$$\Delta u''_m \rightarrow \Delta u'' \text{ weakly in } L^p(0, T; L^2(\Omega)) \quad 1 \leq p < \infty, \quad (44)$$

$$\Delta u''_m \rightarrow \Delta u'' \text{ weakly - star in } L^\infty(0, T; L^2(\Omega)). \quad (45)$$

Thus, by (43) and (44) we get

$$M(a(u_m(t)))\Delta u''_m(t) \rightarrow M(a(u(t)))\Delta u''(t) \text{ weakly in } L^p(0, T; L^2(\Omega)), \quad 1 \leq p < \infty. \quad (46)$$

Since that  $F$  is Lipschitzian, the convergence in (41) implies that

$$F(u_m(t)) \rightarrow F(u(t)) \text{ strongly in } L^p(0, T; L^2(\Omega)). \quad (47)$$

The established convergences make it easy to pass to the limit in the approximate problem (13) and obtain the existence of a strong solution to problem (1).

To prove the continuity established in the Theorem 1, we observe that (from Lemma 2)

$$u_m \in C(0, T; D(A^\alpha)). \quad (48)$$

and

$$u'_m \in C(0, T; D(A^\alpha)). \quad (49)$$

Consequently,

$$u_m \in C^1(0, T; D(A^\alpha)). \quad (50)$$

In order to check up that  $u(0) = u_0$ , it is sufficient to use (48). Analogously, by using (49), it can be shown that  $u'(0) = u_1$ .

**remark 4.1** *If  $M'$  satisfies **(H.1)** then it is possible to show that  $u''' \in L^2(0, T; H^{-1}(\Omega))$ . Consequently, by interpolation, we have that  $u'' \in C([0, T]; L^2(\Omega))$  and, thus the equation  $Ku + F(u) = f$  is satisfied a.e. in  $t$ .*

## 5 Uniqueness

We prove the uniqueness in Theorem 1 as follows.

Let  $u, v$  be solutions of (5) in the class of Theorem 1. It follows that  $w = u - v$  is a solution of:

$$w'' - \Delta w - M(a(u(t)))\Delta w'' - [M(a(v(t))) - M(a(u(t)))] \Delta v'' + F(u(t)) - F(v(t)) = 0, \quad (51)$$

$$w(0) = 0, \quad (52)$$

$$w'(0) = 0. \quad (53)$$

We shall prove that  $w = 0$  in  $[0, T]$ . In fact, taking the inner product in  $L^2(\Omega)$  of (50) by  $2w' \in L^2(0, T; H_0^1(\Omega))$ , we get

$$\begin{aligned} \frac{d}{dt} \left[ |w'(t)|^2 + |\nabla w(t)|^2 + M(a(u(t)))a(w'(t)) \right] &= \frac{d}{dt} (M(a(u(t))))a(w'(t)) \\ &+ [M(a(v(t))) - M(a(u(t)))] 2a(v''(t), w'(t)) \\ &- 2(F(u(t)) - F(v(t)), w'(t)) \end{aligned} \quad (54)$$

i.e.,

$$\begin{aligned}
& \frac{d}{dt} \left[ |w'(t)|^2 + |\nabla w(t)|^2 + M(a(u(t)))a(w'(t)) \right] \quad (55) \\
& \leq \left| \frac{d}{dt} (M(a(u(t)))) \right| |a(w'(t))| \\
& \quad + |M(a(v(t))) - M(a(u(t)))| 2 |a(v''(t), w'(t))| \\
& \quad + |F(u(t)) - F(v(t))| |w'(t)|.
\end{aligned}$$

Observe also that

$$|M(a(v(t))) - M(a(u(t)))| = |M'(\xi)| \left| \|v(t)\|^2 - \|u(t)\|^2 \right|$$

where

$$\xi = (1 - \theta) \|u(t)\|^2 + \theta \|v(t)\|^2, \quad 0 \leq \theta \leq 1.$$

Hence,

$$\begin{aligned}
|M'(\xi)| \left| \|u(t)\|^2 - \|v(t)\|^2 \right| &= |M'(\xi)| (|\|u(t)\| - \|v(t)\||)(\|u(t)\| + \|v(t)\|) \\
&\leq C_2 |\nabla w(t)| (\|u(t)\| + \|v(t)\|)
\end{aligned}$$

because

$$|M'(\xi)| \leq C_2 \text{ if } \xi \in (a(u(t)), a(v(t))) \text{ with } t \in [0, T].$$

Then,

$$\begin{aligned}
& 2 |M(a(u(t))) - M(a(v(t)))| |a(v''(t), w'(t))| \quad (56) \\
& \leq 2C_2 |\nabla w(t)| (\|u(t)\| + \|v(t)\|) |\Delta v''(t)| |w'(t)|.
\end{aligned}$$

Using (6), (8) and the Poincaré's inequality in (55) we obtain that

$$\begin{aligned}
& |M(a(v(t))) - M(a(u(t)))| 2 |a(v''(t), w'(t))| \quad (57) \\
& \leq C_9 \left( |\nabla w(t)|^2 + |\nabla w'(t)|^2 \right).
\end{aligned}$$

Also we have that

$$\begin{aligned}
\left| \frac{d}{dt} M(a(u(t))) \right| &= |M'(a(u(t)))| 2 |a(u(t), u'(t))| \quad (58) \\
&\leq C_2 |\Delta u(t)| |u'(t)| \leq C_{10};
\end{aligned}$$

hence

$$\left| \frac{d}{dt} M(a(u(t))) \right| a(w'(t)) \leq C_{10} a(w'(t)). \quad (59)$$

In the last term of inequality (54), we observe that

$$\begin{aligned} |2(F(u(t)) - F(v(t)), w'(t))| &\leq 2 |F(u(t)) - F(v(t))| |w'(t)| \\ &\leq 2C |w(t)| |w'(t)| \\ &\leq C_{11} [|\nabla w(t)|^2 + |\nabla w'(t)|^2]. \end{aligned} \quad (60)$$

Using (56), (58) and (59) in (54) we obtain

$$\begin{aligned} &\frac{d}{dt} [ |w'(t)|^2 + |\nabla w(t)|^2 + M(a(u(t))) a(w'(t)) ] \\ &\leq C_{12} [ |\nabla w(t)|^2 + |\nabla w'(t)|^2 ] \end{aligned} \quad (61)$$

where

$$C_{12} = \max \{ C_9, C_{10}, C_{11} \} > 0.$$

Integrating (60) from 0 to  $t$  and applying the hypothesis **(A.1)** we have

$$\begin{aligned} &|w'(t)|^2 + |\nabla w(t)|^2 + \rho |\nabla w'(t)|^2 \\ &\leq C_{12} \int_0^t [ |\nabla w(s)|^2 + |\nabla w'(s)|^2 ] ds; \end{aligned} \quad (62)$$

hence by Gronwall inequality we conclude

$$w = 0 \text{ in } [0, T].$$

Thus, the proof of Theorem 1 is completed.  $\square$

**REMARK 2.** In the forthcoming work we will try to study the equation (5) when  $M(\lambda)$  has zero points, that is, degenerate case and  $F(s)$  is a continuous function such that  $sF(s) \geq 0$  for all  $s \in R$ .

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