

ORBITAL STABILITY OF PERIODIC WAVES FOR THE KLEIN-GORDON TYPE EQUATIONS

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This is a joint work with E. Cardoso Jr. - UEM.

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$$E'(\Phi) - cF'(\Phi) = 0. \quad (2)$$

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where $u : \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{C}$ is a complex valued function. We assume that u is an L -periodic function, that is, $u(x + L, t) = u(x, t)$ for all $x, t \in \mathbb{R}$.

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$$F(U) = \text{Im} \int_0^L \bar{u} u_t \, dx = \int_0^L (\text{Re } u \text{ Im } u_t - \text{Im } u \text{ Re } u_t) \, dx.$$

An important qualitative aspect regarding the Hamiltonian systems (1) is the orbital stability.

DEFINITION OF STABILITY.

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$$\sup_{-\infty < t < \infty} \inf\{\|U(t) - \mathbf{T}_1(s_1)\mathbf{T}_2(s_2)\Phi\|_X, -\infty < s_1, s_2 < \infty\} < \varepsilon.$$

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Otherwise, we say that the periodic wave is orbitally unstable in X (or X -unstable).

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(3) $d''(c) > 0$, for all $c \in \mathcal{I}$.

In order to deduce the orbital stability of periodic waves, we need to determine (at least) existence and uniqueness of (weak) solutions.

EXISTENCE AND UNIQUENESS OF WEAK SOLUTIONS

We suppose that $f : \mathbb{R} \rightarrow \mathbb{R}$ is a real function and the Cauchy problem

$$\begin{cases} u_{tt} - u_{xx} + u - f(|u|^2)u = 0 \\ u(x, 0) = u_0(x), \quad u_t(x, 0) = u'_0(x), \end{cases} \quad (7)$$

has a unique (local) solution

$$u \in C([0, T]; H^1_{per}([0, L])) \cap C^1([0, T]; L^2_{per}([0, L])).$$

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In addition, we assume that problem (7) has two (convenient) conserved quantities E and F .

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Theorem 1.

Consider $(u_0, u'_0) \in H_{per}^1([0, L]) \times L_{per}^2([0, L])$. The Cauchy Problem (7) has a unique (local) weak solution $u \in C([0, T]; H_{per}^1([0, L])) \cap C^1([0, T]; L_{per}^2([0, L]))$. In addition, we have the following conserved quantities

$$E(U) = \frac{1}{2} \int_0^L [|u_x|^2 + |u_t|^2 + |u|^2 (2 - \log(|u|^2))] dx,$$

and

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- In addition, we need to use the logarithmic Sobolev inequality

$$4\pi^2 \int_0^L |u|^2 \log |u| \, dx \leq \int_0^L |u_x|^2 \, dx + 2\pi^2 \int_0^L |u|^2 \log \left(\frac{2\pi}{L} \int_0^L |u|^2 \, dx \right) \, dx,$$

$$u \in H_{per}^1([0, L]).$$

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$$\begin{cases} \varphi_{tt} - \varphi_{xx} + \varphi - u \log(|u|^2) + v \log(|v|^2) = 0 \\ \varphi(x, 0) = 0, \quad \varphi_t(x, 0) = 0. \end{cases} \quad (8)$$

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$$\beta(t) \leq \beta_0 e^{-\alpha t},$$

a.e. $t \in [0, T^*]$.

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In addition, the same approach determines sufficient conditions to obtain the spectral property associated with the operator \mathcal{L}_1 .

In order to find a local minimum, we define the orbit generated by Φ

$$\mathcal{O}_{\phi_c} := \left\{ e^{i\theta}(\phi_c(\cdot + y), ic\phi_c(\cdot + y)); (y, \theta) \in [0, L] \times [0, 2\pi) \right\}.$$

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Taking $v := u_t$, let us consider $(y, \theta) \in [0, L] \times [0, 2\pi)$.

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Taking $v := u_t$, let us consider $(y, \theta) \in [0, L] \times [0, 2\pi)$. Let $t \in [0, T]$ be arbitrary but fixed. We define the continuous function

$$\begin{aligned} \Omega_t(y, \theta) &:= \|u_x(\cdot + y, t)e^{i\theta} - \phi'_c\|_{L^2_{per}}^2 \\ &+ (1 - c^2)\|u(\cdot + y, t)e^{i\theta} - \phi_c\|_{L^2_{per}}^2 \\ &+ \|v(\cdot + y, t)e^{i\theta} - ic\phi_c\|_{L^2_{per}}^2. \end{aligned} \quad (11)$$

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Since Ω_t , $t \in [0, T]$, is continuous and $[0, L] \times [0, 2\pi)$ is bounded, we can write,

$$\Omega_t(y(t), \theta(t)) = \inf_{(y, \theta) \in [0, L] \times [0, 2\pi)} \Omega_t(y, \theta) := [\rho_c(\vec{u}(\cdot, t), \mathcal{O}_{\phi_c})]^2. \quad (12)$$

Furthermore, the map

$$t \mapsto \inf_{(y, \theta) \in [0, L] \times [0, 2\pi)} \Omega_t(y, \theta)$$

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$$u(x+y, t)e^{i\theta} := \phi_c(x) + w(x, t) \quad \text{where } w := A + iB \quad (13)$$

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Denoting

$$\vec{w} = (w, z) = (\operatorname{Re} w, \operatorname{Im} z, \operatorname{Im} w, \operatorname{Re} z) = (A, D, B, C).$$

By using the minimum property (above) one has

$$\left\langle \begin{pmatrix} A(\cdot, t) \\ D(\cdot, t) \end{pmatrix}, \begin{pmatrix} \log(\phi_c^2)\phi_c' + 2\phi_c' \\ c\phi_c' \end{pmatrix} \right\rangle_{2,2} = 0 \quad (15)$$

and

$$\left\langle \begin{pmatrix} B(\cdot, t) \\ C(\cdot, t) \end{pmatrix}, \begin{pmatrix} \phi_c \log(\phi_c^2) \\ -c\phi_c \end{pmatrix} \right\rangle_{2,2} = 0, \quad (16)$$

$\forall t \in [0, T]$.

Next, since $G = E - cF$ is a conserved quantity and $G'(\phi_c, ic\phi_c) = (E' - cF')(\phi_c, ic\phi_c) = 0$, we deduce from Taylor's Theorem

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$$\begin{aligned}
 \Delta G &:= G(u_0, u_1) - G(\phi_c, ic\phi_c) \\
 &= G(w(\cdot, t) + \phi_c, z(\cdot, t) + ic\phi_c) - G(\phi_c, ic\phi_c) \\
 &\geq \frac{1}{2} \left\langle \mathcal{L}_R \begin{pmatrix} A(\cdot, t) \\ D(\cdot, t) \end{pmatrix}, \begin{pmatrix} A(\cdot, t) \\ D(\cdot, t) \end{pmatrix} \right\rangle_{2,2} \\
 &\quad + \frac{1}{2} \left\langle \mathcal{L}_I \begin{pmatrix} B(\cdot, t) \\ C(\cdot, t) \end{pmatrix}, \begin{pmatrix} B(\cdot, t) \\ C(\cdot, t) \end{pmatrix} \right\rangle_{2,2} \\
 &\quad - \beta_3 \|\vec{w}(\cdot, t)\|^3 - \beta_4 \|\vec{w}(\cdot, t)\|^4 - \mathcal{O}(\|\vec{w}(\cdot, t)\|^5),
 \end{aligned}$$

Here, operators \mathcal{L}_R and \mathcal{L}_I are defined as

$$\mathcal{L}_R = \begin{pmatrix} -\partial_x^2 + 1 - \log(|\phi_c|^2) - 2 & -c \\ -c & 1 \end{pmatrix} \quad (17)$$

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- Since zero is a simple eigenvalue of \mathcal{L}_1 and $n^-(\mathcal{L}_1) = 1$, we can use the min-max Theorem to guarantee that zero is a simple eigenvalue of \mathcal{L}_R whose eigenfunction is $(\phi'_c, c\phi'_c)$.

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- Since zero is a simple eigenvalue of \mathcal{L}_1 and $n^-(\mathcal{L}_1) = 1$, we can use the min-max Theorem to guarantee that zero is a simple eigenvalue of \mathcal{L}_R whose eigenfunction is $(\phi'_c, c\phi'_c)$. In addition, the min-max Theorem give us that $n^-(\mathcal{L}_R) = 1$.
- The fact that ϕ_c is positive enable us to conclude that zero is the first eigenvalue of \mathcal{L}_I which is simple.

Thus, classical methods of orbital stability (in the sense of definition above) is established on the set

$$\mathcal{A} = \{(u, v) \in H_{per}^1 \times L_{per}^2; F(u, v) = F(\phi_c, ic\phi_c)\},$$

provided that

$$\left\langle \mathcal{L}_{R, \phi_c}^{-1} \begin{pmatrix} c\phi_c \\ \phi_c \end{pmatrix}, \begin{pmatrix} c\phi_c \\ \phi_c \end{pmatrix} \right\rangle_{2,2} = \underbrace{\left\langle \begin{pmatrix} M \\ N \end{pmatrix}, \begin{pmatrix} c\phi_c \\ \phi_c \end{pmatrix} \right\rangle_{2,2}}_{:= -d''(c)} < 0,$$

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where

$$\begin{pmatrix} M \\ N \end{pmatrix} = \begin{pmatrix} \frac{d}{dc}(\phi_c) \\ \phi_c + c \frac{d}{dc}(\phi_c) \end{pmatrix}.$$

However, one has

$$-d''(c) = \int_0^L \phi_c^2 dx + c \underbrace{\frac{d}{dc} \left(\int_0^L \phi_c^2 dx \right)}_{I_c}. \quad (19)$$

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Collecting the results in (19) and (21) we deduce

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that is, $-d''(c) < 0$ if, and only if, $|c| > \frac{\sqrt{2}}{2}$.

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A simple application of the triangle inequality and the fact that G is C^1 map in a neighborhood of the point $(\phi_c, ic\phi_c)$ give us the orbital stability if $(u, v) \notin \mathcal{A}$

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